News of kaons from the lattice

Tom Blum (UConn/RBRC)

FPCP 2022, Ole Miss May 23, 2022

outline

- Direct CP violation (ϵ') in $K \to \pi\pi$ decays (main part of talk)
- Electromagnetic corrections to (semi-) leptonic π / kaon decays
- Kaon mass difference Δm_{κ}
- Long distance contribution to indirect *CP* violation
- Rare decay $K^+ \rightarrow \pi^+ \nu \bar{\nu}$

Direct CP violation in the SM (kaons)

• Final $\pi\pi$ states can have I = 0 or 2.

$$\langle \pi \pi (I=2) | H_w | K^0 \rangle = A_2 e^{i\delta_2} \qquad \Delta I = 3/2 \langle \pi \pi (I=0) | H_w | K^0 \rangle = A_0 e^{i\delta_0} \qquad \Delta I = 1/2$$

- CP symmetry requires A₀ and A₂ be real.
- Direct CP violation in this decay is characterized by:

$$\epsilon' = \frac{i e^{\delta_2 - \delta_0}}{\sqrt{2}} \left| \frac{A_2}{A_0} \right| \left(\frac{\mathrm{Im}A_2}{\mathrm{Re}A_2} - \frac{\mathrm{Im}A_0}{\mathrm{Re}A_0} \right) \quad \begin{array}{c} \text{Direct CP} \\ \text{violation} \end{array}$$

(slide: Norman Christ)

RF2 Snowmass - 5/17/2022 (8)

Direct CP violation in the SM (kaons)

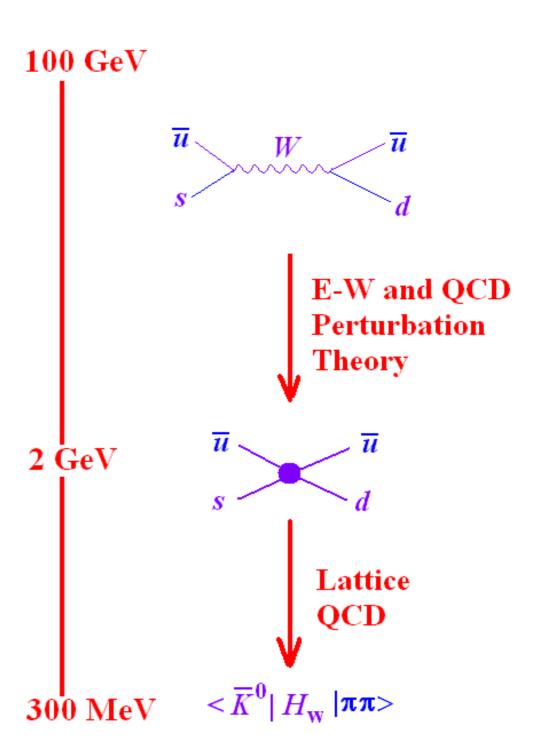
 Represent weak interactions by local four-quark Lagrangian

$$\mathcal{H}^{\Delta S=1} = \frac{G_F}{\sqrt{2}} V_{ud} V_{us}^* \left\{ \sum_{i=1}^{10} \left[z_i(\mu) + \tau y_i(\mu) \right] Q_i \right\}$$

•
$$\tau = -\frac{V_{td}V_{ts}^*}{V_{ud}V_{us}^*} = (1.543 + 0.635i) \times 10^{-3}$$

- $V_{qq'}$ CKM matrix elements
- z_i and y_i Wilson Coefficients
- Q_i four-quark operators

(slide: Norman Christ)



Direct CP violation in the SM (ϵ')

PHYSICAL REVIEW D 102, 054509 (2020)

Editors' Suggestion

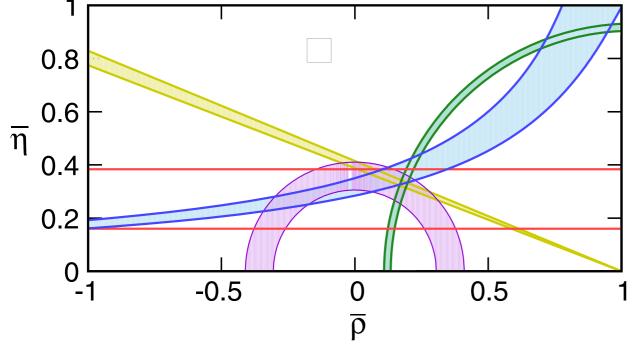
Featured in Physics

Direct *CP* violation and the $\Delta I = 1/2$ rule in $K \rightarrow \pi \pi$ decay from the standard model

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(RBC and UKQCD Collaborations)

$$\begin{aligned} & \operatorname{Re}(\varepsilon'/\varepsilon)_{\mathrm{expt}} = 0.00166(23). & (115) \quad \text{[NA48 2002, KTeV 2011]} \\ & \operatorname{Re}(\varepsilon'/\varepsilon)_{\mathrm{lattice}} = 0.00217(26)(62)(50). & (114) \end{aligned}$$



- G-parity boundary conditions (pions anti-periodic)
- Multiple $\pi\pi$ operators for control of excited states

Errors in the lattice calculation

 $\text{Re}(\varepsilon'/\varepsilon)_{\text{lattice}} = 0.00217(26)(62)(50).$ (114)

statistical, systematic, systematic IB

TABLE XXV. Relative systematic errors on the infinite-volume matrix elements of the $\overline{\text{MS}}$ -renormalized four-quark operators Q'_j .

Error source	Value
Excited state	
Unphysical kinematics	5%
Finite lattice spacing	12%
Lellouch-Lüscher factor	1.5%
Finite-volume corrections	7%
Missing G_1 operator	3%
Renormalization	4%
Total	15.7%

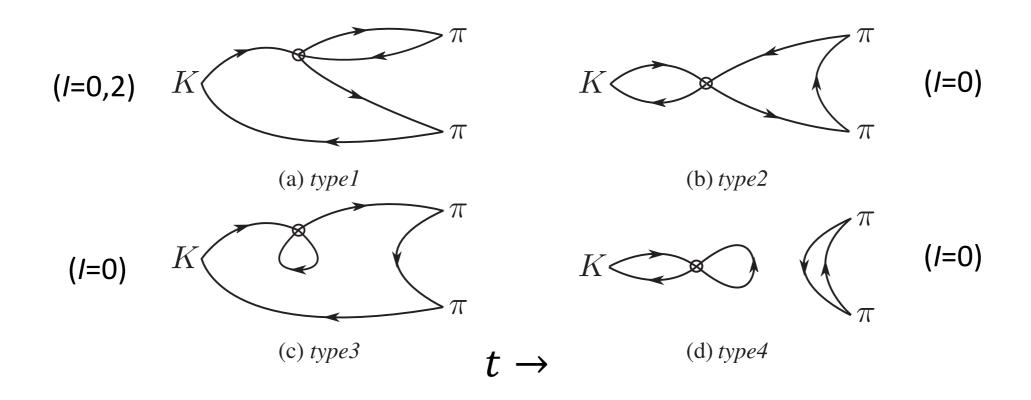
TABLE XXVI.	Relative	systematic	errors	on	$\operatorname{Re}(A_0)$	anc
$\operatorname{Im}(A_0).$						

Error source	Value			
	$\operatorname{Re}(A_0)$	$\operatorname{Im}(A_0)$		
Matrix elements	15.7%	15.7%		
Parametric errors	0.3%	6%		
Wilson coefficients	12%	12%		
Total	19.8%	20.7%		

- Statistical (12%), systematic (29%), IB corrections (23%) [Cirigliano, et al., 2020]
- Continuum PT at NNLO needed for Wilson coefficients (12%)
- New G-parity BC simulations at smaller lattice spacing underway
- EM/isospin corrections can be calculated in lattice QCD+QED, but incompatible with G-parity BC's

Periodic Boundary Conditions (PBCs)

• Every lattice calculation starts with a correlation function in Euclidean spacetime



- C(t) is sum of exponentials, $e^{-E_n t}$, pull operators apart, $t_K \ll t_{op} \ll t_{\pi\pi}$ to get ground state contribution
- G-parity BC's: enforces non-zero momenta for the pions, so 2π ground state corresponds to physical kinematics
- For PBC's an excited state corresponds to physical kinematics

Why make life (seemingly) harder?

- G-parity more expensive ($\sim 2 \times$) and need special ensembles
- G-parity conserves isospin: incompatible with QED corrections
- Ordinary PBC ensembles of gauge configurations exist already with many different lattice spacings, volumes
- Generalized eigenvalue problem (GEVP) method effective for resolving ground and excited states

GEVP method

(Blossier, et al., JHEP 04 (2009) 094)

• Matrix of correlation functions, N operators, N states

 $C_{ij}(t) = \langle O_i(t)O_j(0) \rangle, \quad O_i = \pi\pi, \quad \pi = \overline{\psi}\gamma_5\psi$

• Solve generalized eigenvalue problem,

 $C(t)v_n(t,t_0) = \lambda_n(t,t_0)C(t_0)v_n(t,t_0) \quad 1 \le n \le N, t_0 \le t,$

• Obtain (eigen-) energies of states, eigenvectors

$$\lambda_n(t, t_0) = e^{-E_n(t-t_0)}, \quad v_n \propto u_n, \quad A_n = \sum_{i=1}^N O_i u_{ni}^*,$$

• Systematic error: *N*+1st state contaminates lower *N* states, energies

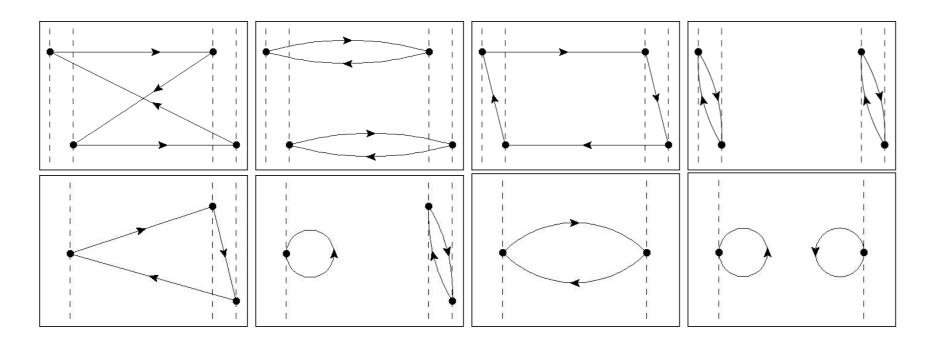
(but the contamination is exponentially suppressed)

$\pi\pi ightarrow \pi\pi$ (Dan Hoying)

• Construct $C_{ij}(t)$, two-pion two-point correlation function

 $C_{ij} = \langle O_i(t)O_j(0) \rangle, \ O_i = \sigma \text{ and } \pi\pi; \ \sigma = \overline{\psi}\psi, \qquad \pi = \overline{\psi}\gamma_5\psi$

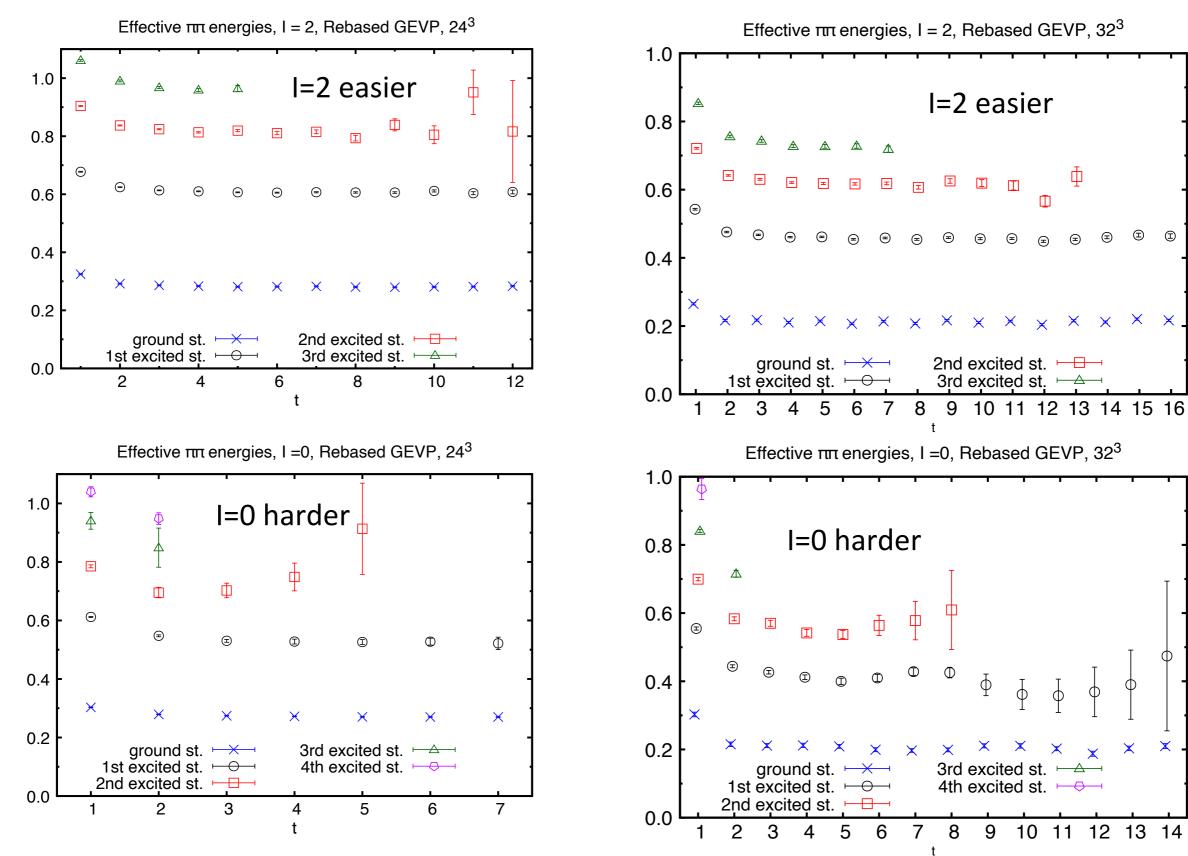
- Use two pion operators with relative momentum up to $\frac{2\pi}{L}(1,1,1)$ and scalar, or *N*=5
- Subtract vacuum state for I=0
- *I*=0, 2 states from linear combinations of Wick contractions



Ensemble details

- 2+1 flavors of Möbius domain-wall fermions (MDWF) with physical masses (m_{π} =140 MeV, *etc*)
- Iwasaki+DSDR gauge action for the gluons
- Two ensembles, 24^3 and 32^3 both with $N_t=64$
- (inverse) lattice spacing $a^{-1} = 1$ and 1.4 GeV (relatively coarse)
- Statistics: 258 and 107 configs for 24³ and 32³, respectively
- 32³ ensemble matches G-parity BC ensemble (741 configs)

$\pi\pi \to \pi\pi$ energies



$\pi\pi \to \pi\pi$ operators n = 1, Evec, Rebased GEVP 1.0 aE_{eff} ⊢⊖−− C(000) - X- $C(\sigma)$ • Elements of eigenvector *n* C(001) ⊢ <u>△</u> 0.5 C(011) ⊢⊟ · are coefficients of linear ¥ Ж C(111) ⊢⊖⊣ Ж Ж ¥ Ж × combination of operators 0.0 E đ đ đ I Π Ш that create state *n* 24³ -0.5 ¥ 4 ⋠ ⋠ 4 -1.0 2 3 5 6 4 t/a preliminary n = 1, Evec, Rebased GEVP 1.0 aE_{eff} ⊢⊖− /=0 C(000) -**32³** $C(\sigma)$ 5x5 GEVP C(001) -0.5 C(011) ⊢⊟ · C(111) Noise from higher states 0.0 "Rebase": use linear đ Г П $\overline{\Phi}$ Ē Ш Ф ф combination from small t for -0.5 reduced GEVP at larger t -1.0 10 12 2 6 8 14 t/a

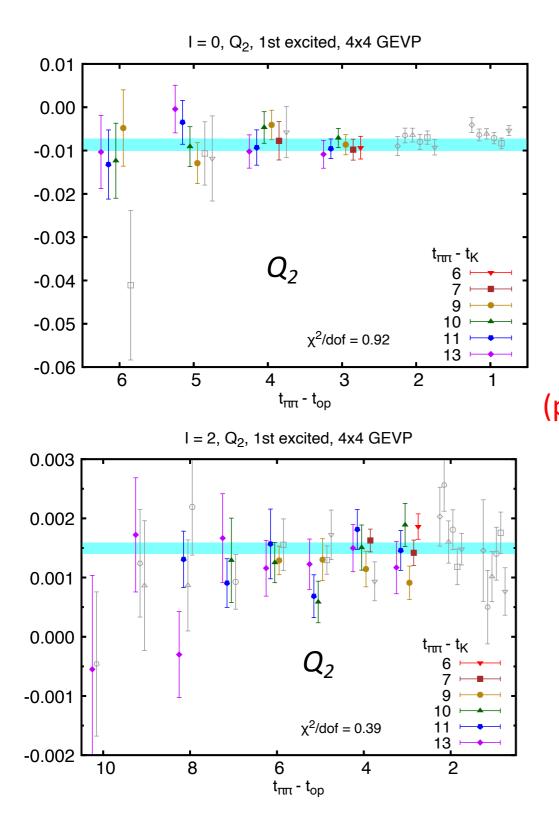
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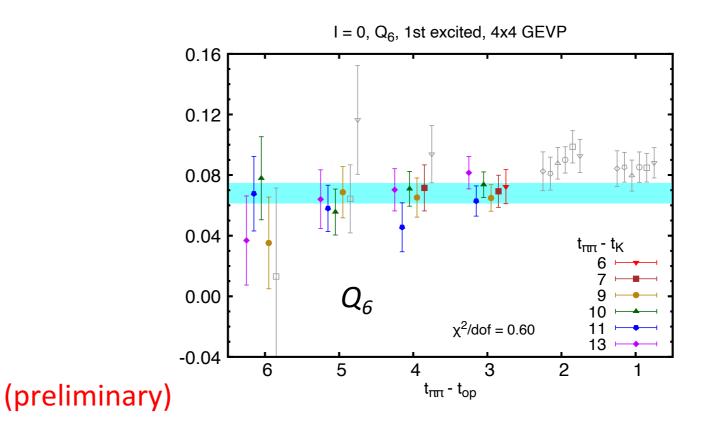
$K ightarrow \pi \pi$ (Masaaki Tomii)

✓ Use 2-pion states from $\pi\pi$ scattering, *i.e.*, construct linear combinations of correlation functions using coefficients from $\pi\pi$ analysis

Could also do GEVP for kaon (under investigation, but not yet)

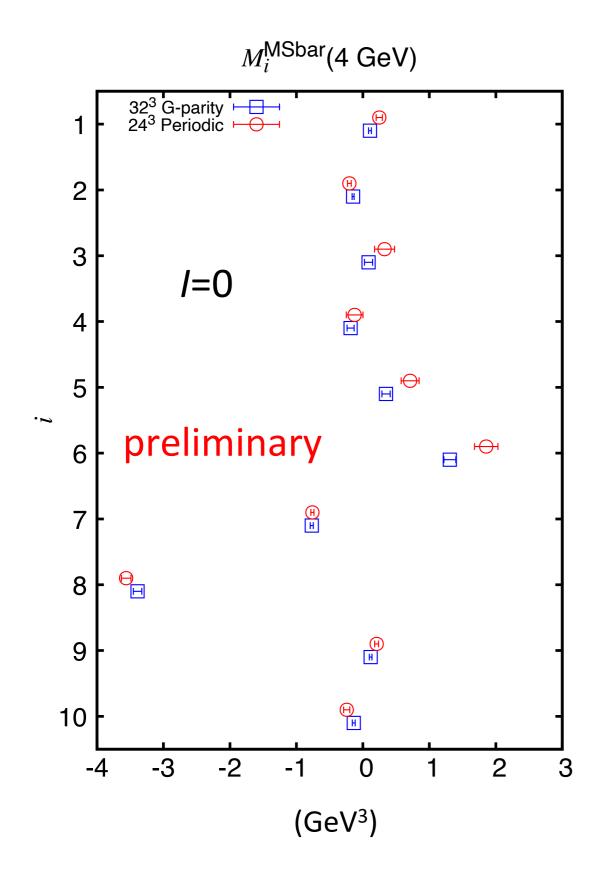
$K \rightarrow \pi \pi$ matrix elements





- Filled symbols used in fit
- Fit to constant
- Several (fixed) $t_{\pi\pi} t_K$ separations
- unrenormalized
- *I*=2 not optimized (sparsening for *I*=0)

Comparison with G-parity BC



- PBC: 24³, *a*⁻¹= 1 GeV
- GPBC: 32³, *a*⁻¹= 1.4 GeV (3x stats)
- Renormalized matrix elements
- Consistency is encouraging
- 32³ in progress
- Results: A_0, A_2, ϵ' at Lattice 2022 in August

Outlook

- GEVP method, with similar statistics, appears promising
- Improving statistics on both ensembles
- Moving to finer lattice spacings (including for G-Parity)
- Method for EM corrections under development– large hurdle cleared with PBC setup
- Apply to next generation of 2+1+1 MDWF ensembles
- 10% precision within next 10 years

Leptonic decay rates from
lattice QCD+QED

$$\Gamma(K^{+} \rightarrow \mu^{+}\nu_{\mu}) \propto (1 + \delta_{EM})|V_{us}|^{2}f_{K^{+}}^{2}$$

$$\frac{\Gamma(K \rightarrow \mu\nu_{\mu})}{\Gamma(\pi \rightarrow \mu\nu_{\mu})} = 1.3367(28) (PDG)$$

$$\frac{|V_{us}|f_{K^{+}}}{|V_{ud}|f_{\pi^{+}}} = 0.27600(37)$$

$$\frac{|V_{us}|f_{K^{+}}}{|V_{ud}|f_{\pi^{+}}} = 0.27683(29)(20)$$

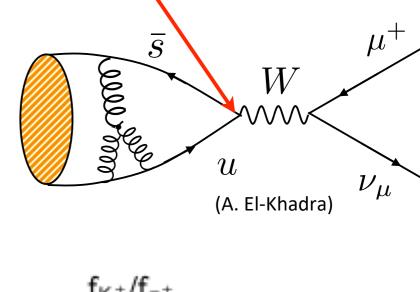
$$\frac{|V_{us}|f_{K^{+}}}{|V_{ud}|f_{\pi^{+}}} = 0.27600(4)$$

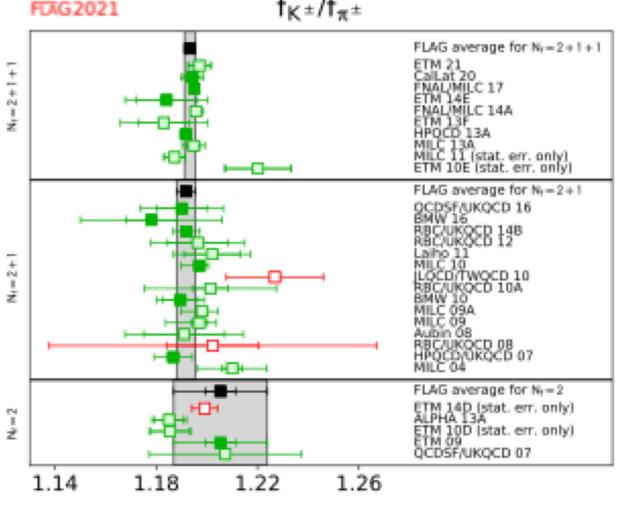
$$\frac{|V_{us}|f_{K^{+}}}{|V_{ud}|f_{\pi^{+}}} = 0.2760(4)$$

$$\frac{|V_{us}|f_{K^{+}}}{|V_{ud}|f_{\pi^{+}}} = 0.2760(4)$$

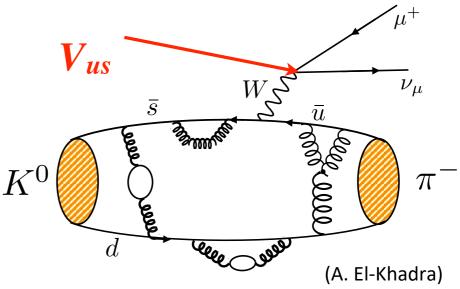
$$\frac{f_{K^{+}}}{|V_{ud}|f_{\pi^{+}}} = 0.2760(4)$$

$$\frac{f_{K^{+}}}{|V_{ud}|f_{\pi^{+}}} = 1.1932(21)$$
0.18%





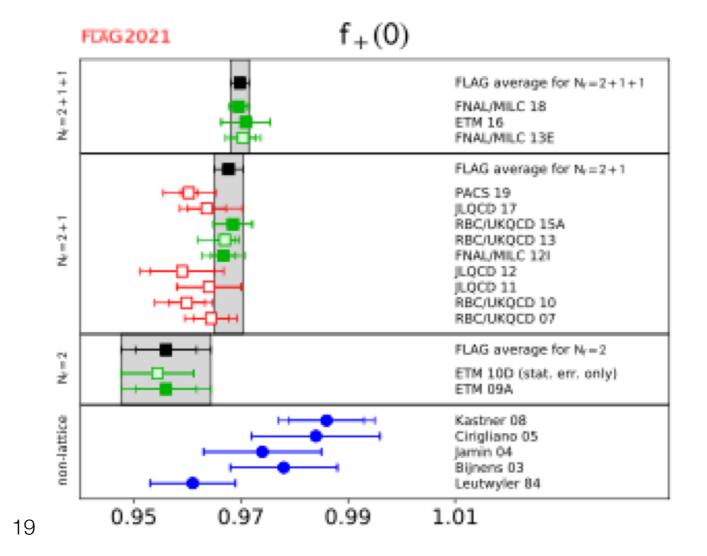
Semi-leptonic decay rates from lattice QCD



 $\Gamma\left(K^0 \to \pi^- \mu^+ \nu_\mu\right) \propto \left(1 + \frac{\delta_{EM}}{\delta_{SU(2)}}\right) |V_{us}|^2 f_+^2(0)$

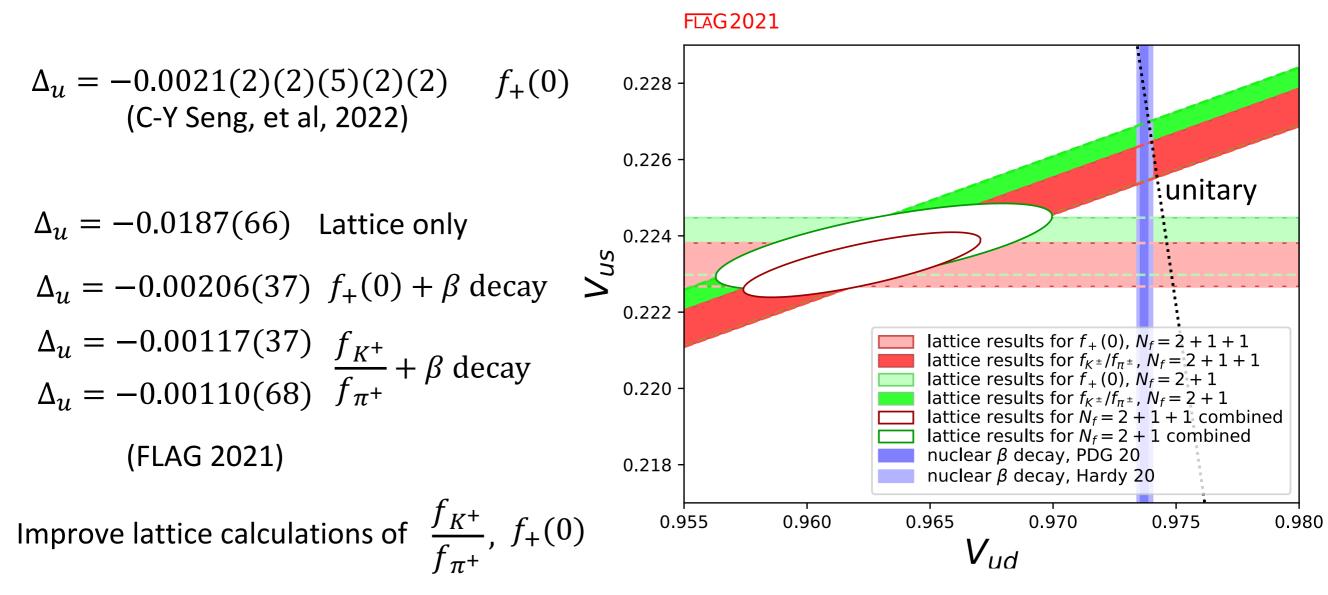
 $|V_{us}|f_+(0) = 0.21635(39)(3)$ (C-Y Seng, et al, 2022)

 $|V_{us}|f_{+}(0) = 0.2165(4)$ $f_{+}(0) = 0.9698(17)$ (FLAG 2021)



Tests of 1st row CKM unitarity

$$\Delta_u \equiv |V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 - 1$$



Light-meson leptonic decay rates in lattice QCD+QED

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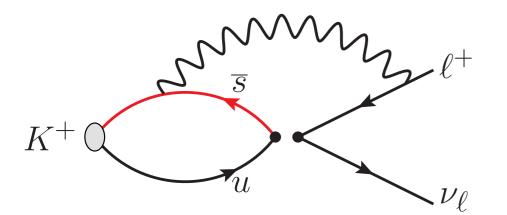
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- First lattice QCD+QED calculation
- 1.53(19) % and 0.24(10) % effects on π, K leptonic decay rates
- V_{us}=0.22538(46) (2x red. over PDG)
- QED_L (power law FV effects)

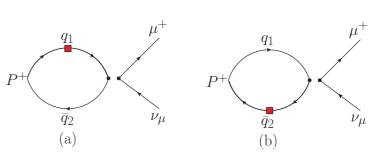


Figure 1: Feynman diagrams of scalar insertions on quark legs (marked with red boxes).

(e)

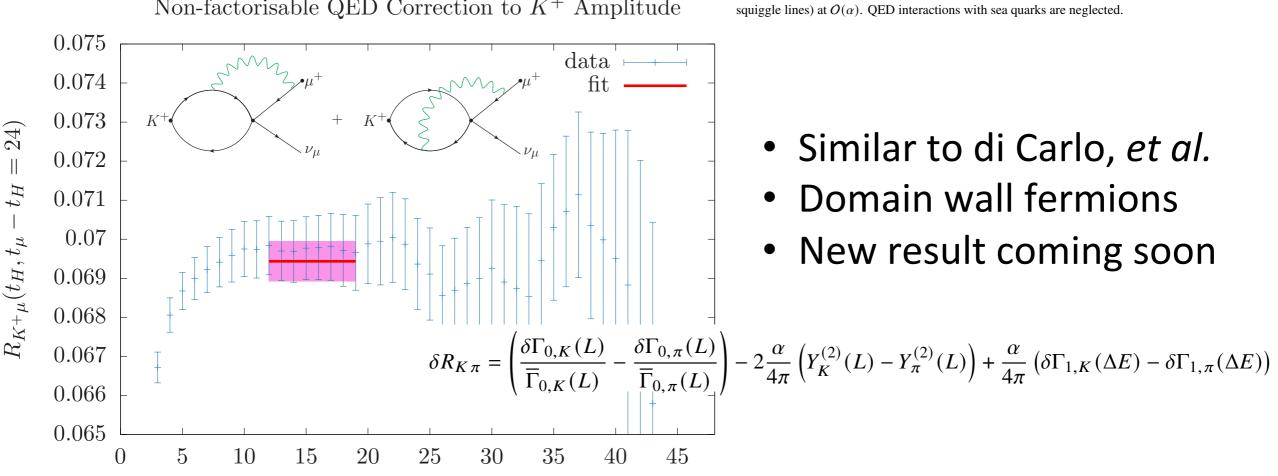
Figure 2: Feynman diagrams of all possible insertions of the electromagnetic current (marked with green

Near-Physical Point Lattice Calculation of Isospin-Breaking Corrections to $K_{\ell 2}/\pi_{\ell 2}$

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Non-factorisable QED Correction to K^+ Amplitude

 t_H



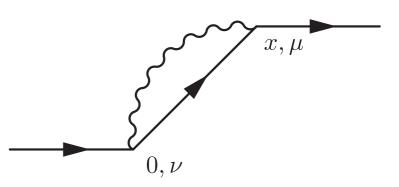
Infinite volume reconstruction for EM corrections

QED self-energies from lattice **QCD** without power-law finite-volume errors

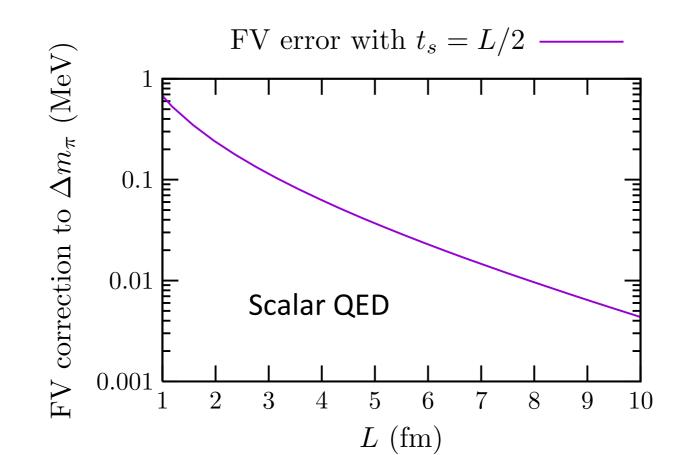
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- Naïve treatment: long-distance part has power-law FVE
- Break up time integral for loop into short and long distance pieces (t_s)
- If excited states suppressed: QCD in finite volume, QED treated infinite volume
- FVE are exponentially suppressed, not power-law



Infinite volume reconstruction: pion mass splitting

Lattice QCD Calculation of the Pion Mass Splitting

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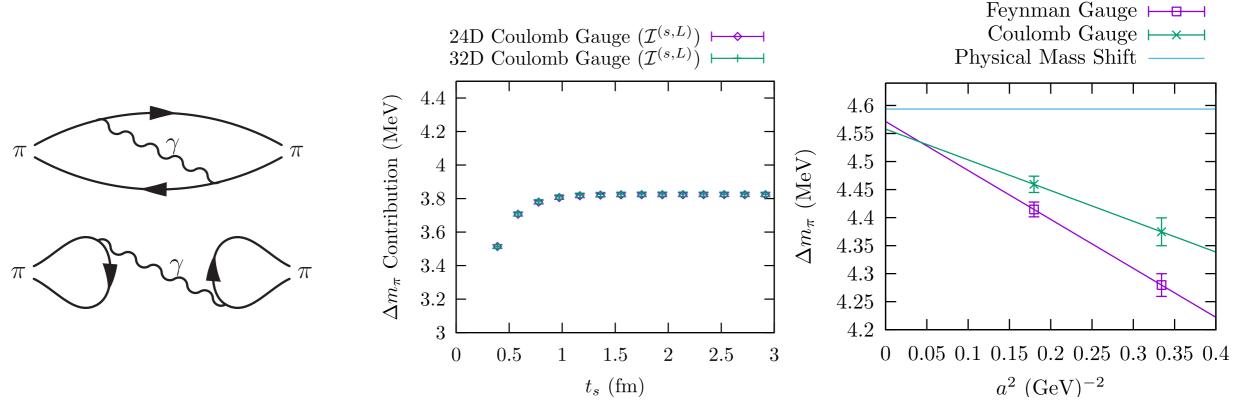


TABLE I. Previous lattice calculations of $m_{\pi^{\pm}} - m_{\pi^{0}}$ are compared to this Letter. Note $m_{\pi^{\pm}}$ is the charged pion mass $m_{\pi^{0}}$ is the neutral pion mass

$m_{\pi^{\pm}} - m_{\pi^0}$ (MeV)
$5.33(48)_{stat}(59)_{sys}^{a}$
$4.60(20)_{\text{stat}}$
$4.21(23)_{\text{stat}}(13)_{\text{sys}}$
$4.534(42)_{stat}(43)_{sys}$

- 5x error reduction
- Pion, kaon leptonic decays soon

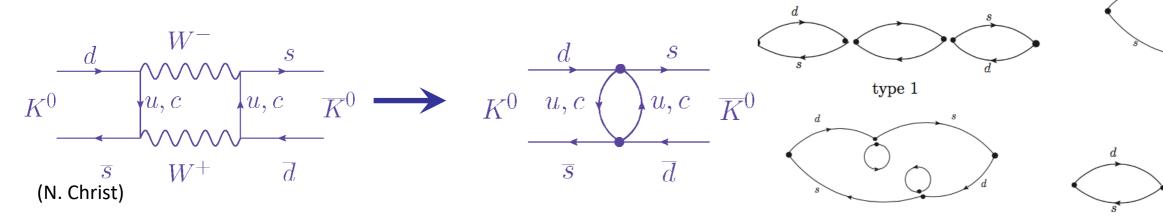
Outlook

✓ 1st row CKM unitarity interesting, needs further attention from lattice community

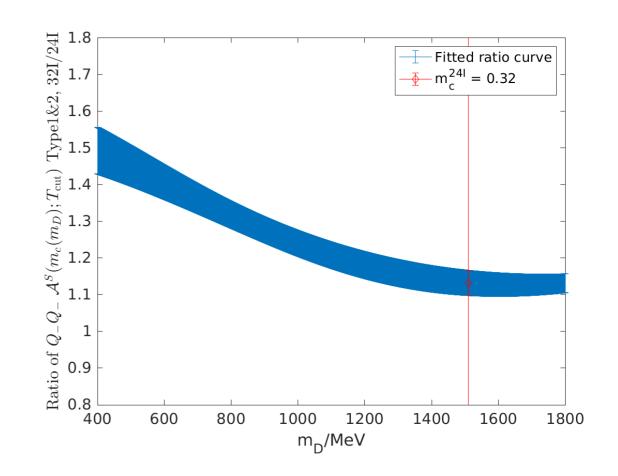
✓ Improvement of lattice calculations of $\frac{f_K}{f_{\pi}}$, $f_+(0)$ from 0.2% precision → sub 0.1 % seems doable

 Treatment of EM effects important at this level of precision, encouraging progress in several directions

$K_L - K_S$ mass difference Δm_K (Bigeng Wang, RBC-UKQCD)



- Exp: 3.482(6) × 10⁻¹² MeV
- 1000 TeV reach
- Lattice: 5.8(0.6)_{stat}(2.3)_{sys}10⁻¹² MeV
- Single lattice, a⁻¹=2.36 GeV
- Lattice: 10% stat errors, 1% in 10 yrs
- 40% non-zero lattice spacing error biggest hurdle (unexpected)
- c-quark important (GIM)
- Dynamical charm, small a: exascale resources needed (coming soon)
- 5% precision goal by 2026



type 3

type 2

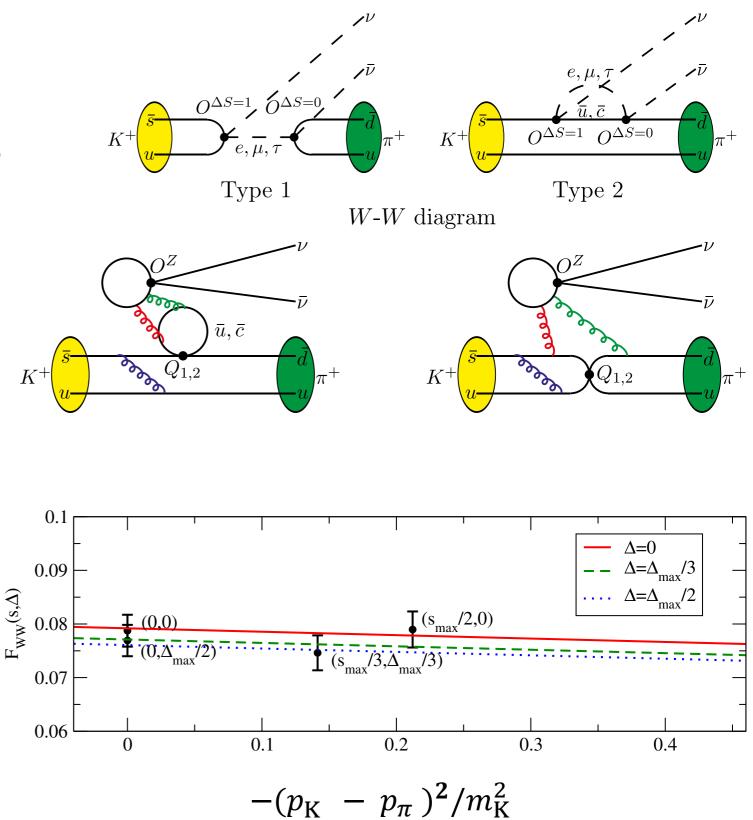
type 4

Indirect CP violation and E

- sub-% precision: need ~6 % long-distance contributions
- Closely related to mixing matrix for Δm_{κ}
- Added complication: log divergence, needs non-perturb. Renormalization
- Subtraction of divergence induces new 4-quark operator with low-energy constant computable in QCD PT
- Similar systematics / goals as for Δm_K

$K^+ \rightarrow \pi^+ \nu \bar{\nu} decays$ [RBC/UKQCD 2016-2019]

- Highly suppressed, good BSM probe
- E949, 2x SM, 60-70% uncertainty
- NA68 10% precision goal
- Theory clean, but long-distance (LD) effects 5-10%
- Developed method for LD effects, completed pilot lattice calculations
- Physical kinematics calculation underway on a⁻¹=2.7 GeV lattice with 30% target precision



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