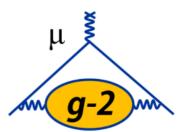






# **Experimental Review of Muon g-2**

Esra Barlas-Yucel on behalf of the Muon g-2 Collaboration FPCP - University of Mississippi 25 May 2022

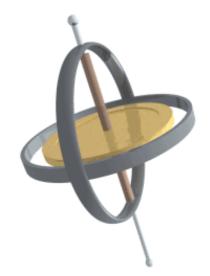




## Magnetic Moment of Muon

$$\overrightarrow{\mu} = g_{\mu} \frac{e}{2m} \overrightarrow{s}$$

g: Proportionality constant between spin and magnetic moment



## **Anomalous Magnetic Moment of Muon**

$$a_{\mu} = \frac{g_{\mu} - 2}{2}, \qquad \overrightarrow{\mu} = (1 + a_{\mu}) \frac{e}{m} \overrightarrow{s}$$

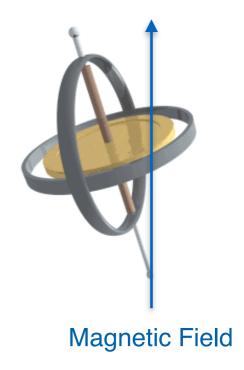
Shows how much g differs fractionally from 2!

Measuring this anomaly could tell us if there are new particles or even forces that contribute to  $a_{\boldsymbol{\mu}}$ 

## Magnetic Moment of Muon

$$\overrightarrow{\mu} = g_{\mu} \frac{e}{2m} \overrightarrow{s}$$

g: Proportionality constant between spin and magnetic moment



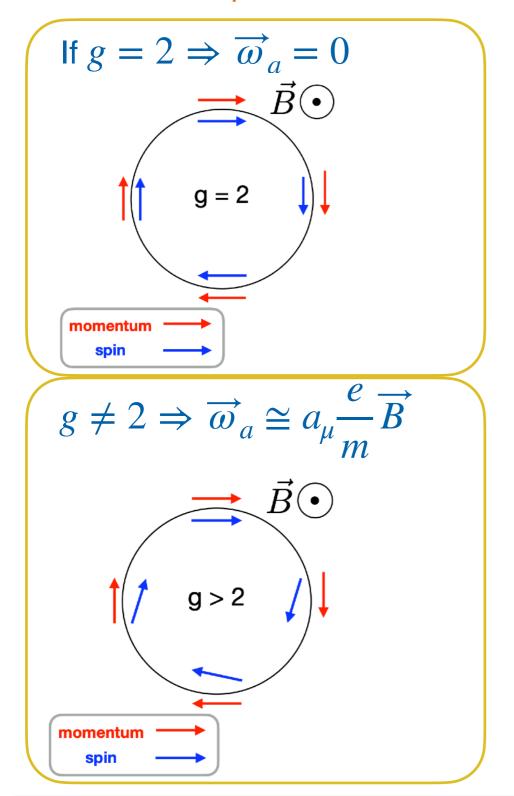
## **Anomalous Magnetic Moment of Muon**

$$a_{\mu} = \frac{g_{\mu} - 2}{2}, \qquad \overrightarrow{\mu} = (1 + a_{\mu}) \frac{e}{m} \overrightarrow{s}$$

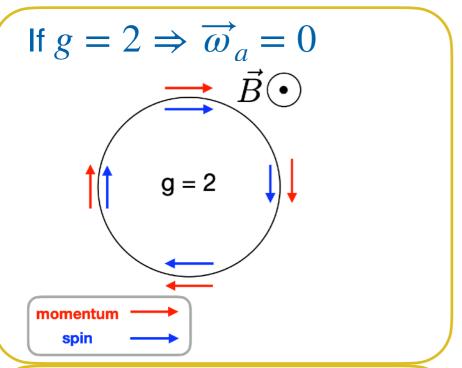
Shows how much g differs fractionally from 2!

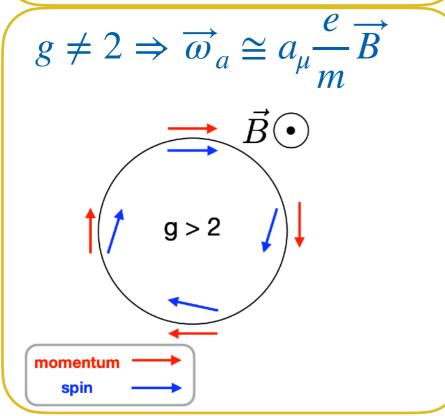
Measuring this anomaly could tell us if there are new particles or even forces that contribute to  $a_{\boldsymbol{\mu}}$ 

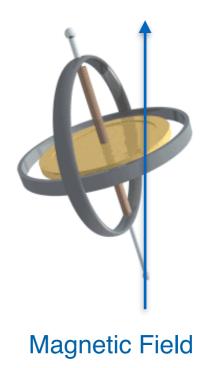
polarized muons in a magnetic field



polarized muons in a magnetic field

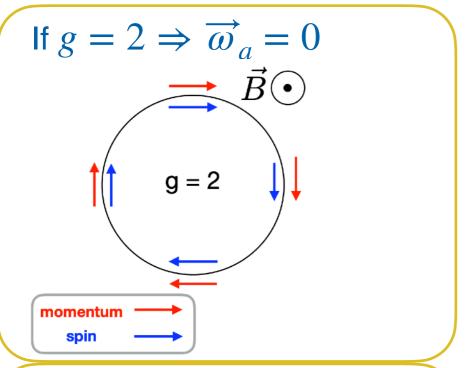


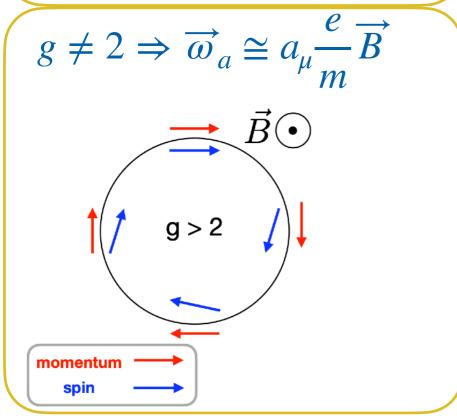


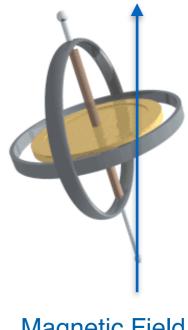


$$\overrightarrow{\omega}_c = -\frac{e}{\gamma m}\overrightarrow{B}$$
 , cyclotron frequency (freq. of charged particle under magnetic field)

#### polarized muons in a magnetic field





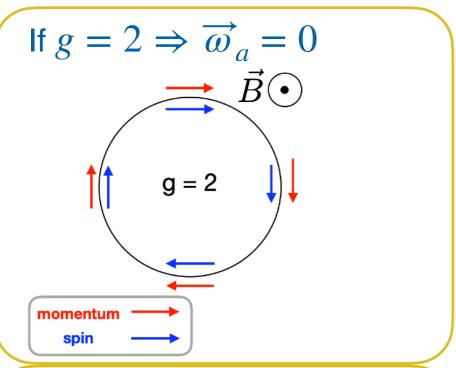


Magnetic Field

$$\overrightarrow{\omega}_c = -\frac{e}{\gamma m}\overrightarrow{B}$$
 , cyclotron frequency (freq. of charged particle under magnetic field)

$$\overrightarrow{\omega}_{\scriptscriptstyle S} = -\, \frac{e}{\gamma m} \overrightarrow{B} (1 + \gamma a_{\mu})$$
 , Larmor precession frequency(total spin precession freq.)

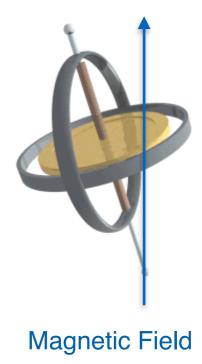
#### polarized muons in a magnetic field



$$g \neq 2 \Rightarrow \overrightarrow{\omega}_a \cong a_\mu \frac{e}{m} \overrightarrow{B}$$

$$g > 2$$

$$momentum \Rightarrow spin$$



$$\overrightarrow{\omega}_c = -\frac{e}{vm}\overrightarrow{B}$$
 , cyclotron frequency (freq. of charged particle under magnetic field)

$$\overrightarrow{\omega}_{s}=-rac{e}{\gamma m}\overrightarrow{B}(1+\gamma a_{\mu})$$
 , Larmor precession frequency(total spin precession freq.)

$$\overrightarrow{\omega}_a\cong\overrightarrow{\omega}_s-\overrightarrow{\omega}_c$$
 , anomalous precession frequency 
$$\overrightarrow{\omega}_a\cong a_\mu\frac{e}{m}\overrightarrow{B}$$

Measure them to extract anomaly

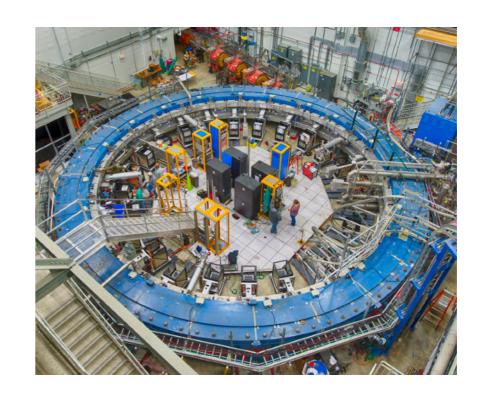
# **CERN and BNL g-2 Experiments**

- Muon g-2 experiments were launched at CERN to measure  $a_{\mu}$ 
  - Launched a measurement (with a long dipole magnet) with precision of 2% and finally **0.4%** (1961,1962)
  - Second measurement with a storage ring-> x25 more precise- **270ppm**-> Revealed a quantitative discrepancy with the theory (1966)
  - Third measurement with a storage ring + electrostatic quadrupole field-> **7.3 ppm** precision (1969-1979)
- BNL took over the effort on measuring the anomaly more precisely (1997-2001)
  - Storage ring magnet + electrostatic quadrupole field + Fast Kickers +24 calorimeters
  - $\mu^+$  and  $\mu^-$  running
  - A factor of 14 improvement on precision after latest CERN measurement **0.54ppm**

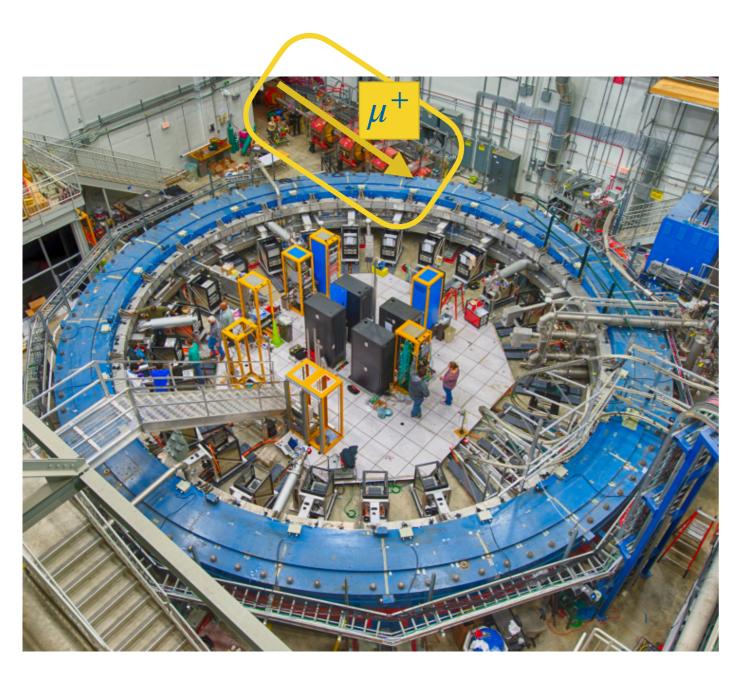
# Fermilab Muon g-2 Experiment





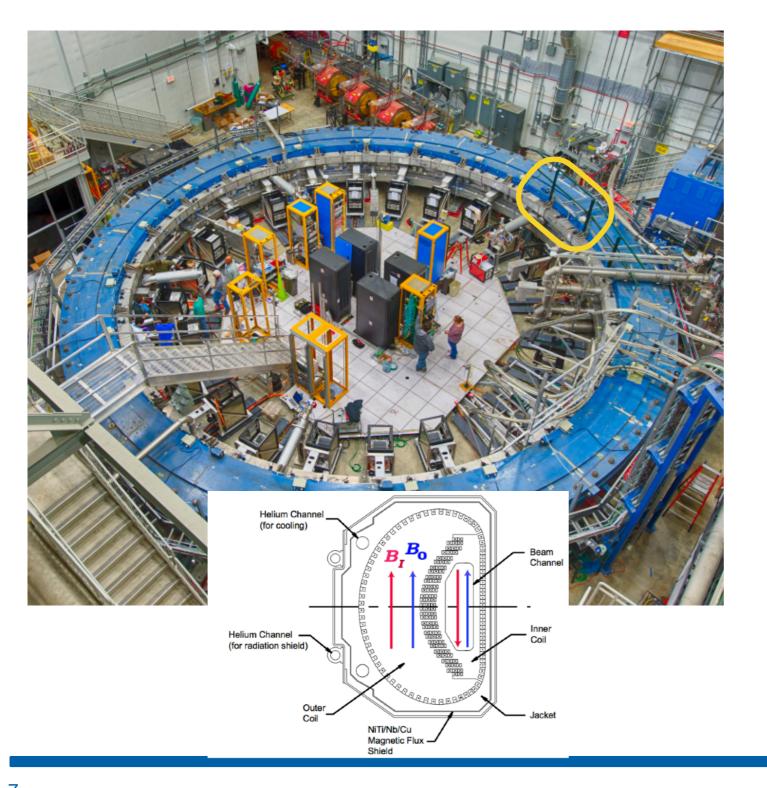


# **Muons at the Experimental Hall!**



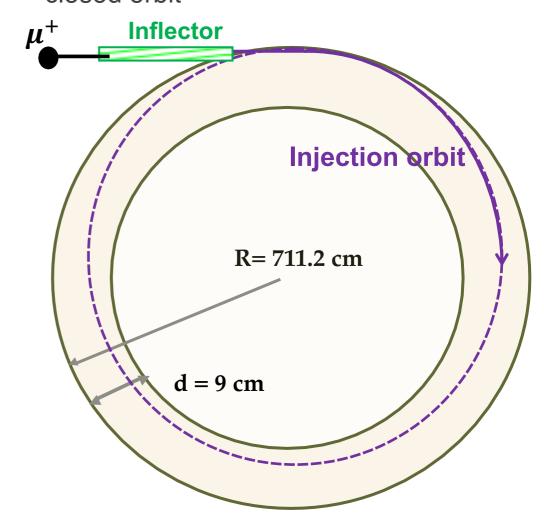
- 8 GeV protons are delivered to Recycler Ring from Booster
- Split the proton bunch into four bunches with RF system.
- Direct the proton punches to pion production target and obtain pions.
- Muons produced by pion decays circulate in the delivery ring until proton contamination is removed.
- Deliver 3.1 GeV polarized muons to g-2 storage ring .

## **Storing the Muons: Inflector and Kickers**

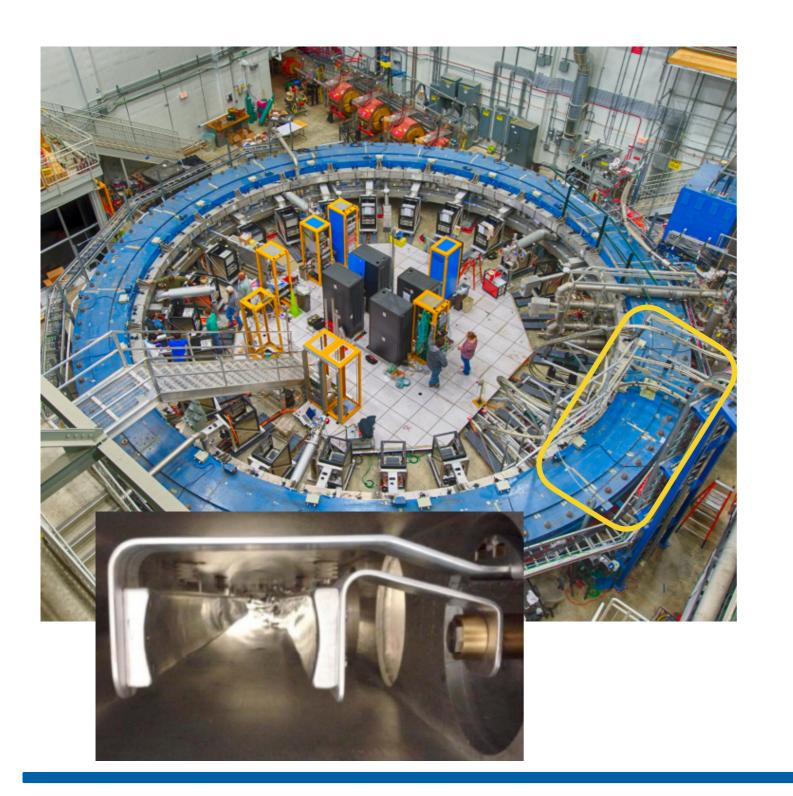


#### Inflector

- Super conducting magnet
- Cancels B field(1.45T) in the magnet gap and let the beam enter the storage ring without being deflected.
- They are at r=77mm outside central closed orbit

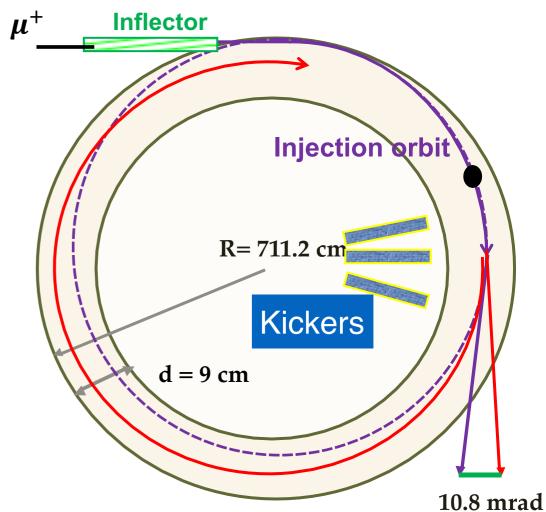


## **Storing the Muons: Inflector and Kickers**



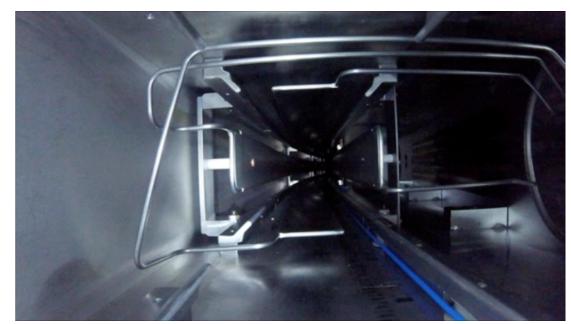
#### Magnetic Kickers

- Kick some more to direct the muons into ideal orbit.
- Use 10.8 mrad pulsed kicks (<149 ns)



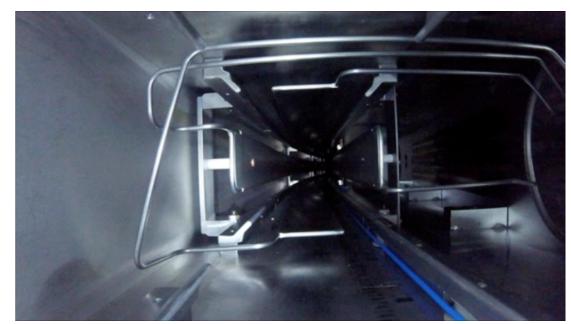


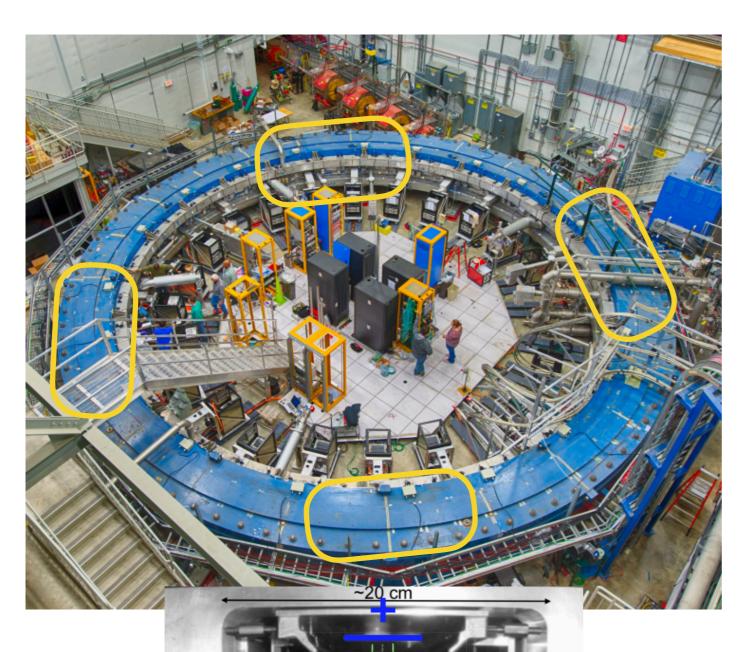
- Electrostatic Quadrupoles
  - 4 sets of quads which cover 43% of the ring
  - Electrostatic quadrupoles are used to focus the beam vertically while the storage ring field provides horizontal focusing
  - Cancels out leading order of electric field contribution running at magic momentum p = 3.094 GeV/c



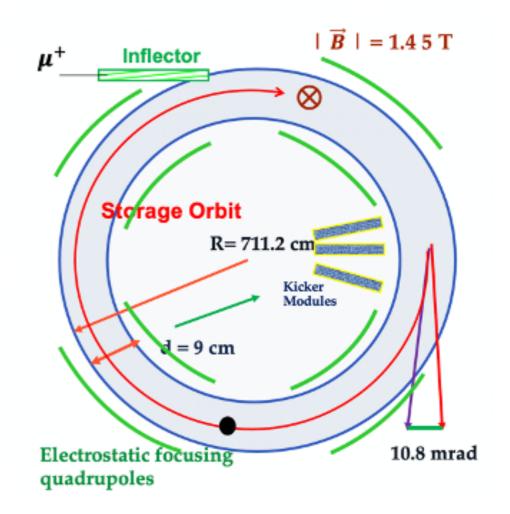


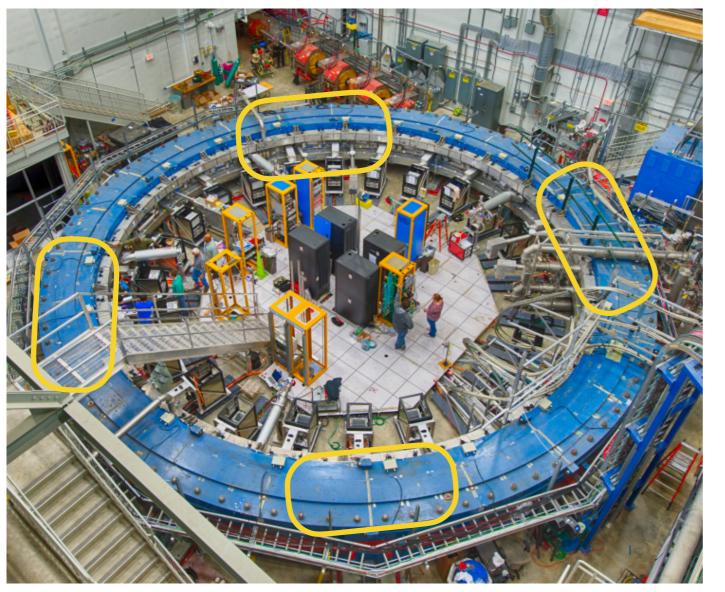
- Electrostatic Quadrupoles
  - 4 sets of quads which cover 43% of the ring
  - Electrostatic quadrupoles are used to focus the beam vertically while the storage ring field provides horizontal focusing
  - Cancels out leading order of electric field contribution running at magic momentum p = 3.094 GeV/c





- Electrostatic Quadrupoles
  - 4 sets of quads which cover 43% of the ring
  - Electrostatic quadrupoles are used to focus the beam vertically while the storage ring field provides horizontal focusing
  - Cancels out leading order of electric field contribution running at magic momentum p = 3.094 GeV/c



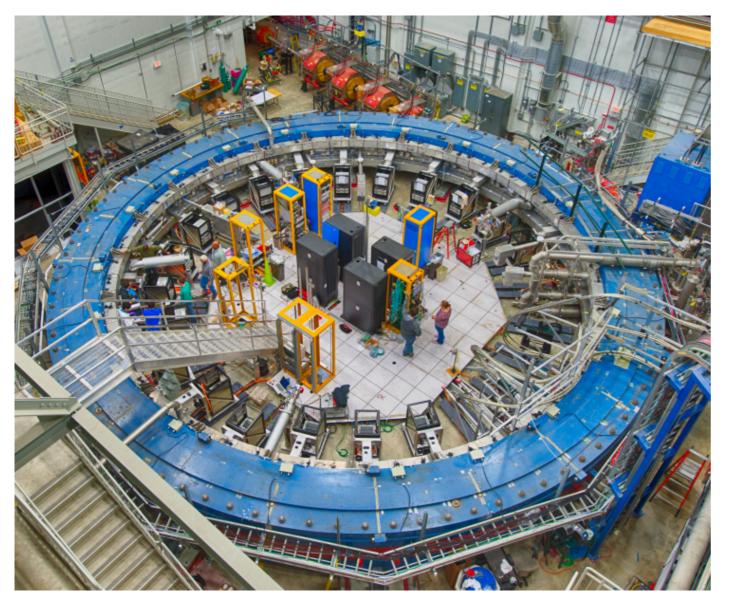


#### Electrostatic Quadrupoles

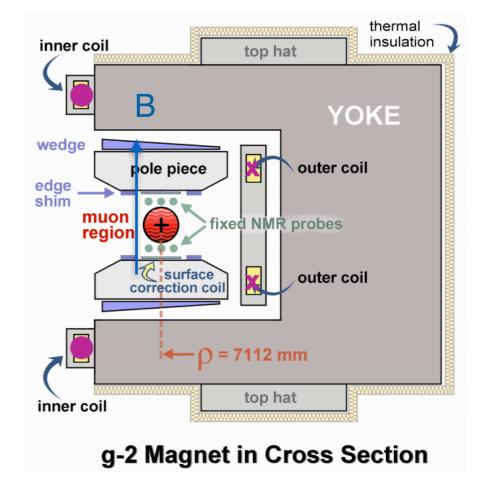
- 4 sets of quads which cover 43% of the ring
- Electrostatic quadrupoles are used to focus the beam vertically while the storage ring field provides horizontal focusing
- Cancels out leading order of electric field contribution running at magic momentum p = 3.094 GeV/c

$$\vec{\omega}_{a} = \frac{e}{m} \left[ a_{\mu} \vec{B} - a_{\mu} \frac{\gamma}{\gamma + 1} (\vec{\beta} \cdot \vec{B}) \vec{\beta} - \left( a_{\mu} - \frac{1}{\gamma^{2} - 1} \right) \vec{\beta} \times \vec{E} \right]$$
0 if in plane
Term cancels at the magic momentum

# **Storing the Muons: Magnet**



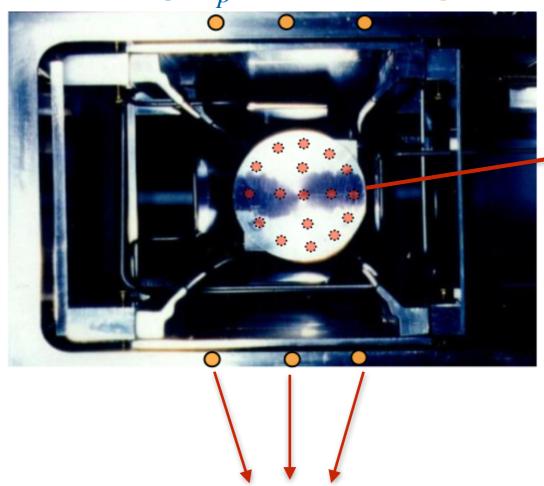
Achieved 25 ppm on field uniformity

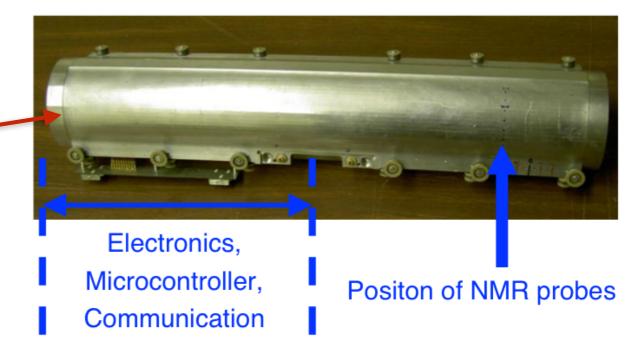


Superconducting C shaped magnet
Provides 1.45T B field(vertical and uniform)

- 12 Yokes: Open on the inside, allows the decay positrons to reach to the detectors.
- 72 poles:Low-carbon steel to minimize the impurity
- 144 Edge shims: Minimize the local sextupole field by changing edge shim thickness
- 864 Steel wedges: Angle adjustment (compensate quadrupole component), radial adjustment (shim local dipole field).
- Surface correction coil: Reduces nonuniformities on higher moment of field.

# Measuring $\omega_p$ : Monitoring and Measuring the Magnetic Field





#### Fixed probes:

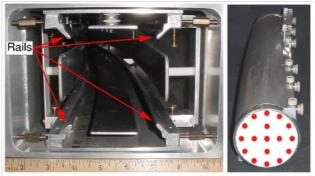
- 378 probes located on vacuum chamber
- Measure the magnetic field while muons are inside the storage ring

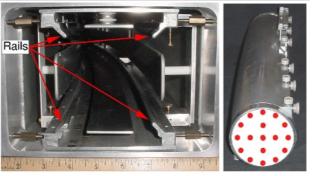
### Trolley(Motorized cart):

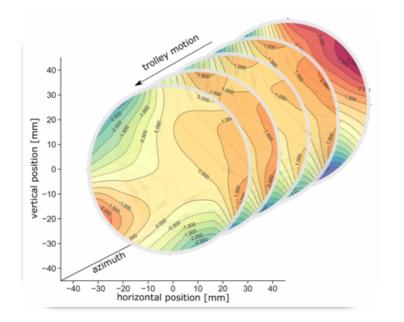
- 17 NMR probes
- Circles around the ring on periodically
- Measures the magnetic field in the storage region
- Used to calibrate FP measurements

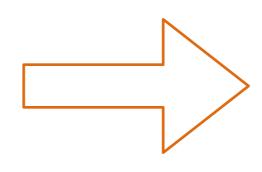
# Measuring $\omega_p$ : Monitoring and Measuring the Magnetic Field



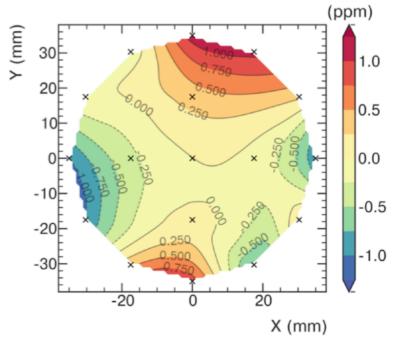








## azimuthally averaged field



## To determine $\omega_p$ at all times:

- Map the magnetic field in the storage region with trolley runs every 3 days
- Use fixed probes to interpolate the field between trolley runs

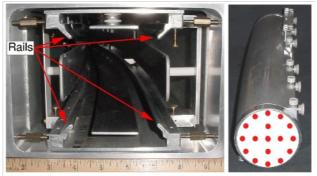
$$a_{\mu} = \left(\frac{g_e}{2}\right) \left(\frac{\omega_a}{\langle \omega_p \rangle}\right) \left(\frac{\mu_p}{\mu_e}\right) \left(\frac{m_{\mu}}{m_e}\right)$$

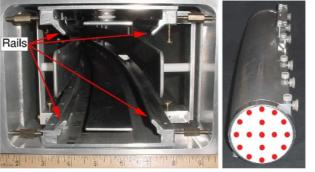
$$\langle \omega_p \rangle \approx \omega_p \otimes \rho(r)$$

 $\omega_{p}$ : free proton precession frequency Using proton NMR  $\hbar\omega_p=2\mu_p B$ 

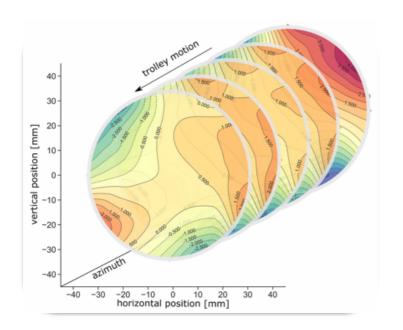
# Measuring $\omega_p$ : Monitoring and Measuring the Magnetic Field

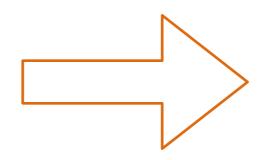


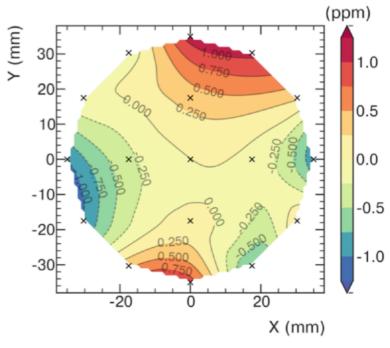




## azimuthally averaged field



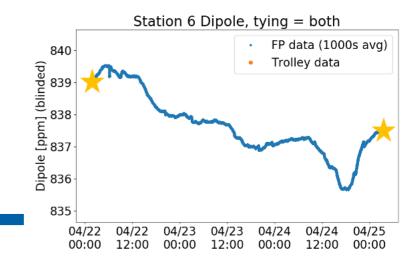




## To determine $\omega_p$ at all times:

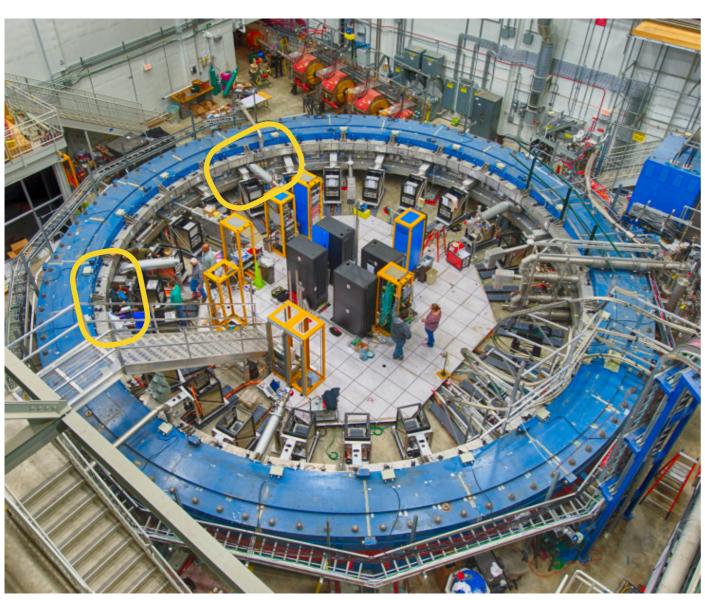
- Map the magnetic field in the storage region with trolley runs every 3 days
- Use fixed probes to interpolate the field between trolley runs

$$a_{\mu} = \left(\frac{g_e}{2}\right) \left(\frac{\omega_a}{\langle \omega_p \rangle}\right) \left(\frac{\mu_p}{\mu_e}\right) \left(\frac{m_{\mu}}{m_e}\right)$$
 $\langle \omega_p \rangle \approx \omega_p \otimes 
ho(r)$ 



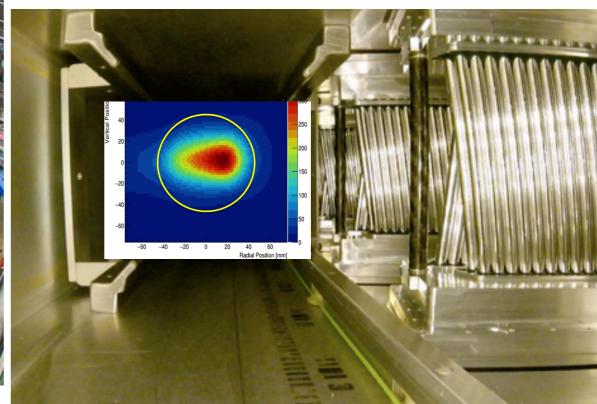
# **Detectors: Trackers for Reconstructing the Beam Profile**





#### Trackers

- 2 straw-tracker stations
- 8 modules per station each with
   128 straws
- Reconstruct muon beam profile from positron trajectories

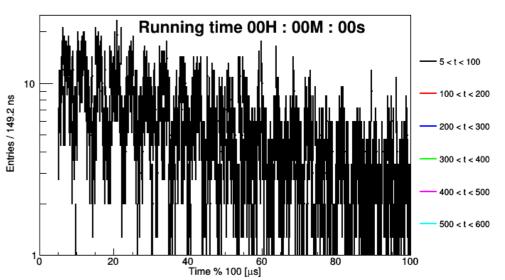


$$a_{\mu} = \left(\frac{g_e}{2}\right) \left(\frac{\omega_a}{\langle \omega_p \rangle}\right) \left(\frac{\mu_p}{\mu_e}\right) \left(\frac{m_{\mu}}{m_e}\right)$$

$$\langle \omega_p \rangle \approx \omega_p \otimes \rho(r)$$

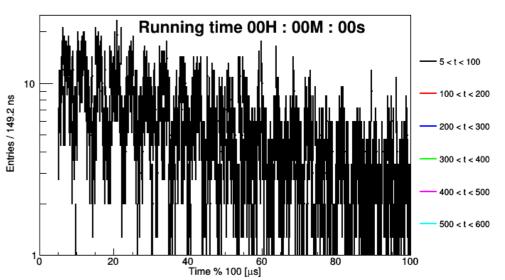


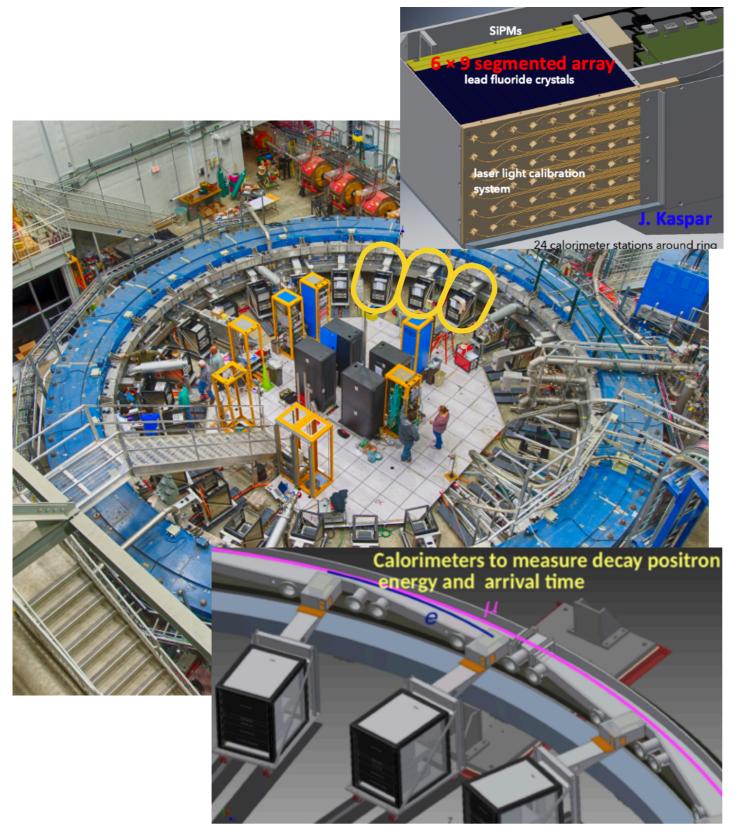
- 24 segmented PbF<sub>2</sub> crystal calorimeters stationed around the ring
- Detects energy and arrival time of  $e^+$  decayed from muons:  $\mu^+ \to e^+ \bar{\nu_\mu} \nu_e$



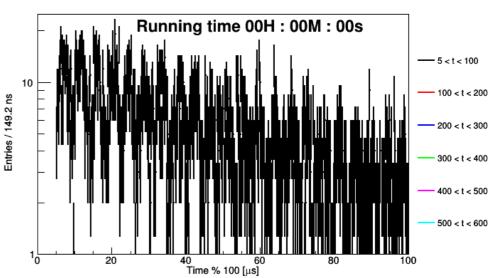


- 24 segmented PbF<sub>2</sub> crystal calorimeters stationed around the ring
- Detects energy and arrival time of  $e^+$  decayed from muons:  $\mu^+ \to e^+ \bar{\nu_\mu} \nu_e$



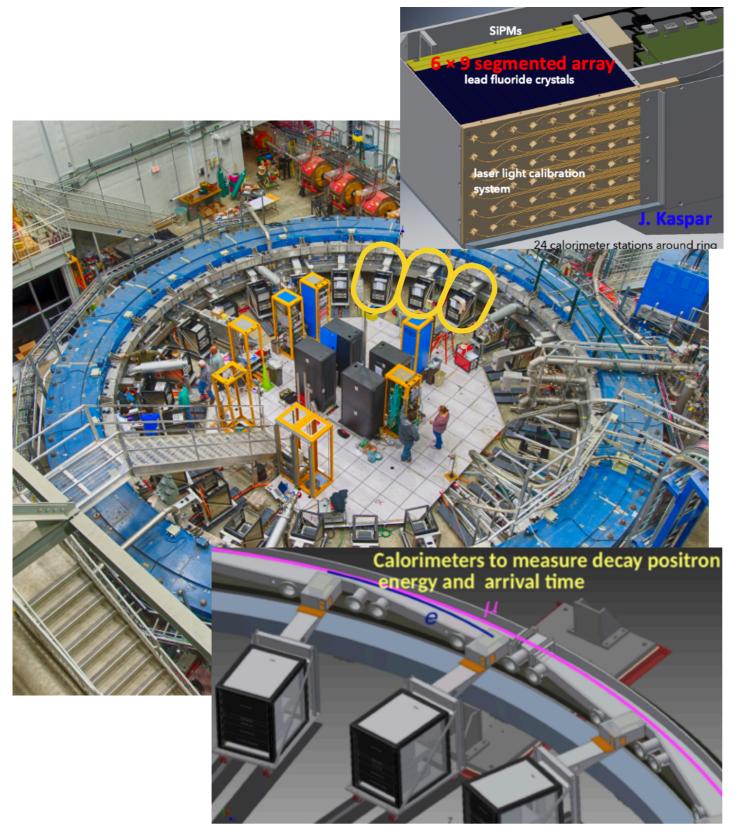


- 24 segmented PbF<sub>2</sub> crystal calorimeters stationed around the ring
- Detects energy and arrival time of  $e^+$  decayed from muons:  $\mu^+ \to e^+ \bar{\nu_{\mu}} \nu_e$

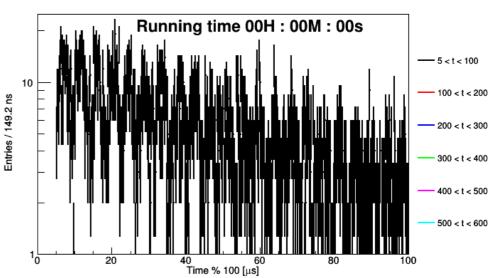


$$N(t) = N_0 e^{-t/\tau} [1 + A\cos(\omega_a t + \phi)]$$

$$a_{\mu} = \left(\frac{g_e}{2}\right) \left(\frac{\omega_a}{\langle \omega_p \rangle}\right) \left(\frac{\mu_p}{\mu_e}\right) \left(\frac{m_{\mu}}{m_e}\right)$$



- 24 segmented PbF<sub>2</sub> crystal calorimeters stationed around the ring
- Detects energy and arrival time of  $e^+$  decayed from muons:  $\mu^+ \to e^+ \bar{\nu_{\mu}} \nu_e$

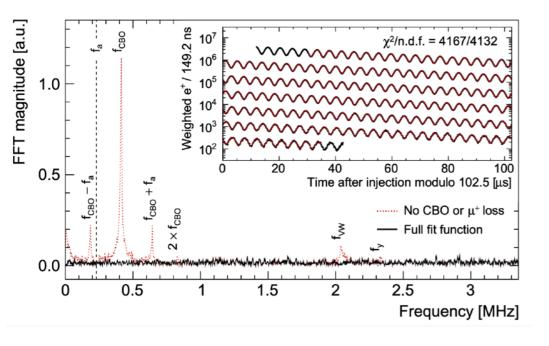


$$N(t) = N_0 e^{-t/\tau} [1 + A\cos(\omega_a t + \phi)]$$

$$a_{\mu} = \left(\frac{g_e}{2}\right) \left(\frac{\omega_a}{\langle \omega_p \rangle}\right) \left(\frac{\mu_p}{\mu_e}\right) \left(\frac{m_{\mu}}{m_e}\right)$$

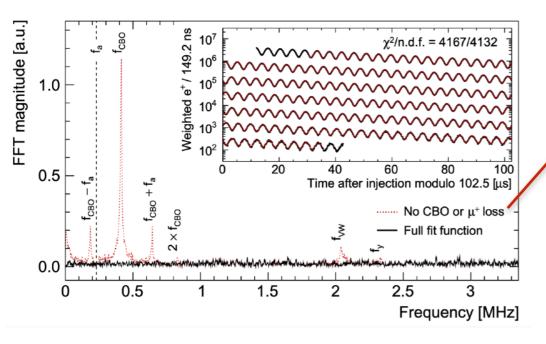


## FFT analysis of fit residuals





### FFT analysis of fit residuals



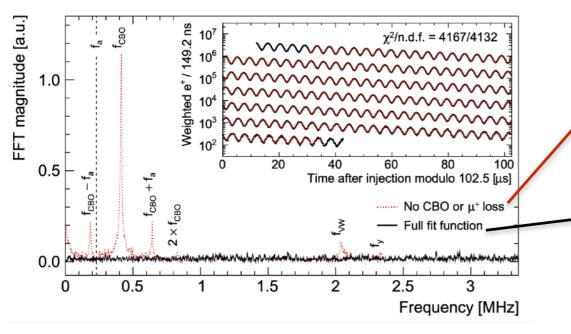
### **Underling Physics**

5 parameter fit function

$$N(t) = N_0 e^{-t/\tau} [1 + A\cos(\omega_a t + \phi)]$$



### FFT analysis of fit residuals



#### **Underling Physics**

#### 5 parameter fit function

$$N(t) = N_0 e^{-t/\tau} [1 + A\cos(\omega_a t + \phi)]$$

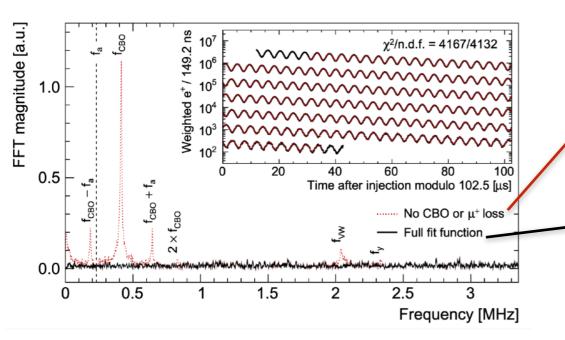
#### **Systematic Effects**

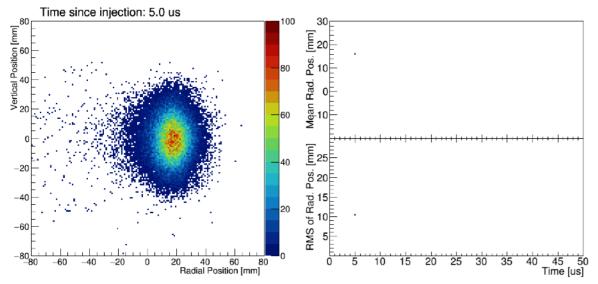
Including CBO, lost muon, other beam dynamics related parameters improve the fit results

$$\begin{split} N_0 \, e^{-\frac{t}{\gamma \tau}} \left( 1 + \boldsymbol{A} \cdot \boldsymbol{A}_{BO}(t) \cos(\omega_a \, t + \phi \cdot \phi_{BO}(t) \,) \, \cdot \, N_{\text{CBO}}(t) \cdot \, N_{\text{VW}}(t) \cdot \, N_y(t) \cdot \, N_{2\text{CBO}}(t) \cdot \, J(t) \right. \\ \left. \boldsymbol{A}_{\text{BO}}(t) = 1 + \boldsymbol{A}_A \cos(\omega_{\text{CBO}}(t) + \phi_A) e^{-\frac{t}{\tau_{\text{CBO}}}} \right. \\ \left. \boldsymbol{\phi}_{\text{BO}}(t) = 1 + \boldsymbol{A}_\phi \cos(\omega_{\text{CBO}}(t) + \phi_\phi) e^{-\frac{t}{\tau_{\text{CBO}}}} \right. \\ \left. \boldsymbol{N}_{\text{CBO}}(t) = 1 + \boldsymbol{A}_{\text{CBO}} \cos(\omega_{\text{CBO}}(t) + \phi_{\text{CBO}}) e^{-\frac{t}{\tau_{\text{CBO}}}} \right. \\ \left. \boldsymbol{N}_{\text{2CBO}}(t) = 1 + \boldsymbol{A}_{\text{2CBO}} \cos(2\omega_{\text{CBO}}(t) + \phi_{\text{2CBO}}) e^{-\frac{t}{2\tau_{\text{CBO}}}} \right. \\ \left. \boldsymbol{N}_{\text{VW}}(t) = 1 + \boldsymbol{A}_{\text{VW}} \cos(\omega_{\text{VW}}(t) t + \phi_{\text{VW}}) e^{-\frac{t}{\tau_{\text{VW}}}} \right. \\ \left. \boldsymbol{N}_y(t) = 1 + \boldsymbol{A}_y \cos(\omega_y(t) t + \phi_y) e^{-\frac{t}{\tau_y}} \right. \\ \left. \boldsymbol{J}(t) = 1 - \boldsymbol{k}_{LM} \int_{t_0}^t \boldsymbol{\Lambda}(t) dt \right. \\ \left. \boldsymbol{\omega}_{\text{CBO}}(t) = \omega_0 t + \boldsymbol{A} e^{-\frac{t}{\tau_A}} + \boldsymbol{B} e^{-\frac{t}{\tau_B}} \right. \\ \left. \boldsymbol{\omega}_y(t) = \boldsymbol{F} \boldsymbol{\omega}_{\text{CBO}(t)} \sqrt{2\omega_c/\boldsymbol{F}\omega_{\text{CBO}}(t) - 1} \right. \\ \left. \boldsymbol{\omega}_{\text{VW}}(t) = \omega_c - 2\omega_y(t) \end{split}$$



## FFT analysis of fit residuals





### **Underling Physics**

5 parameter fit function

$$N(t) = N_0 e^{-t/\tau} [1 + A\cos(\omega_a t + \phi)]$$

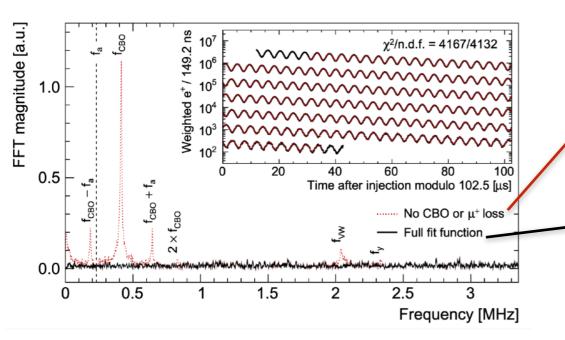
#### **Systematic Effects**

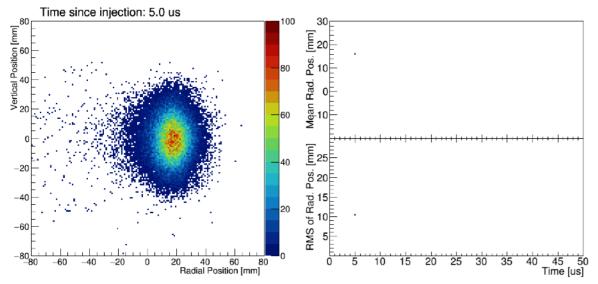
Including CBO, lost muon, other beam dynamics related parameters improve the fit results

$$\begin{split} N_0 \, e^{-\frac{t}{\gamma \gamma}} \left( 1 + A \cdot A_{BO}(t) \cos(\omega_a \, t + \phi \cdot \phi_{BO}(t) \,) \, \cdot \, N_{\text{CBO}}(t) \cdot \, N_{\text{VW}}(t) \cdot \, N_y(t) \cdot \, N_{2\text{CBO}}(t) \cdot \, J(t) \right. \\ \left. A_{\text{BO}}(t) = 1 + A_A \cos(\omega_{\text{CBO}}(t) + \phi_A) e^{-\frac{t}{\gamma_{\text{CBO}}}} \right. \\ \left. \phi_{\text{BO}}(t) = 1 + A_\phi \cos(\omega_{\text{CBO}}(t) + \phi_\phi) e^{-\frac{t}{\gamma_{\text{CBO}}}} \right. \\ \left. N_{\text{CBO}}(t) = 1 + A_{\text{CBO}} \cos(\omega_{\text{CBO}}(t) + \phi_{\text{CBO}}) e^{-\frac{t}{\gamma_{\text{CBO}}}} \right. \\ \left. N_{\text{2CBO}}(t) = 1 + A_{2\text{CBO}} \cos(2\omega_{\text{CBO}}(t) + \phi_{2\text{CBO}}) e^{-\frac{t}{\gamma_{\text{CBO}}}} \right. \\ \left. N_{\text{VW}}(t) = 1 + A_{\text{VW}} \cos(\omega_{\text{VW}}(t) t + \phi_{\text{VW}}) e^{-\frac{t}{\gamma_{\text{VW}}}} \right. \\ \left. N_y(t) = 1 + A_y \cos(\omega_y(t) t + \phi_y) e^{-\frac{t}{\gamma_y}} \right. \\ \left. J(t) = 1 - k_{LM} \int_{t_0}^t \Lambda(t) dt \right. \\ \left. \omega_{\text{CBO}}(t) = \omega_0 t + A e^{-\frac{t}{\gamma_A}} + B e^{-\frac{t}{\gamma_B}} \right. \\ \left. \omega_y(t) = F \omega_{\text{CBO}(t)} \sqrt{2\omega_c/F} \omega_{\text{CBO}}(t) - 1 \right. \\ \left. \omega_{\text{VW}}(t) = \omega_c - 2\omega_y(t) \end{split}$$



## FFT analysis of fit residuals





### **Underling Physics**

5 parameter fit function

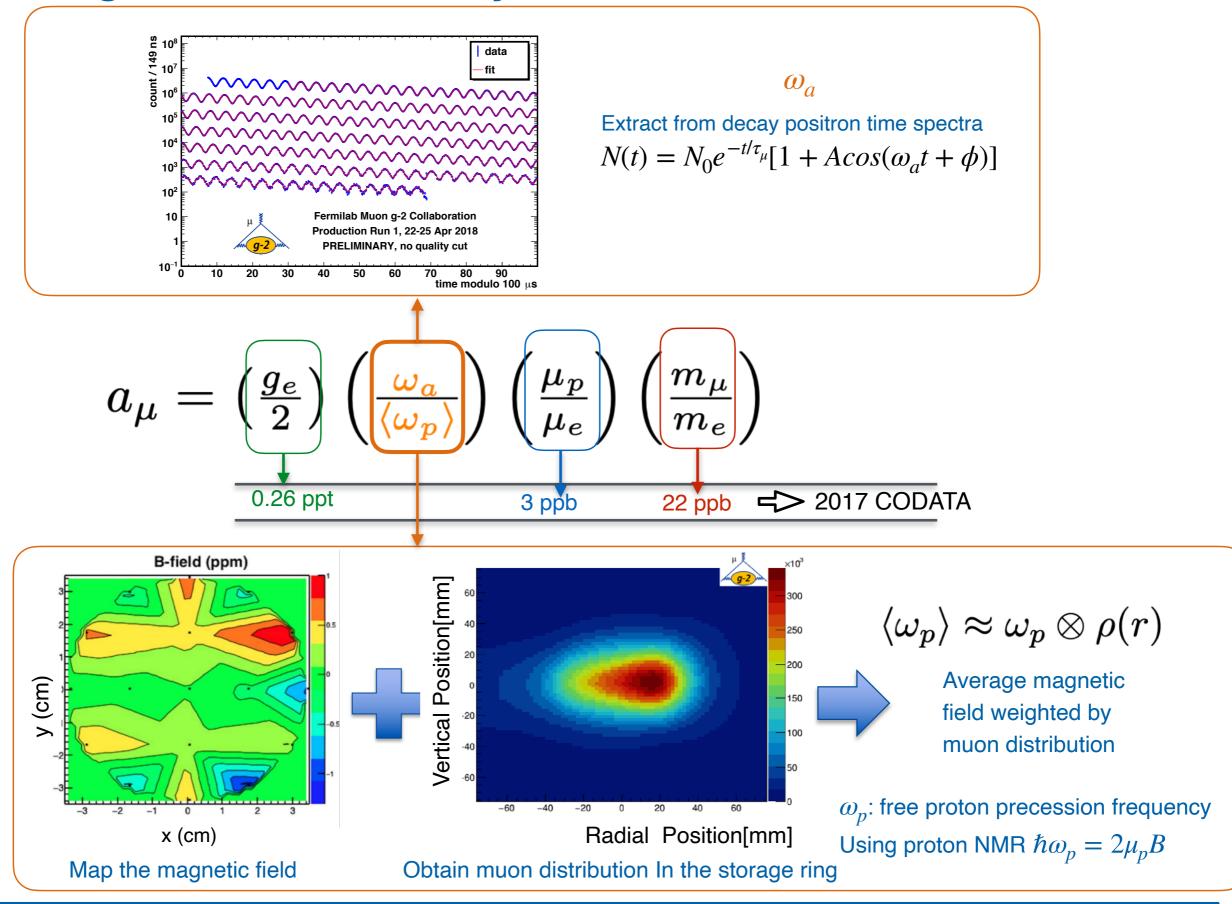
$$N(t) = N_0 e^{-t/\tau} [1 + A\cos(\omega_a t + \phi)]$$

#### **Systematic Effects**

Including CBO, lost muon, other beam dynamics related parameters improve the fit results

$$\begin{split} N_0 \, e^{-\frac{t}{\gamma \gamma}} \left( 1 + A \cdot A_{BO}(t) \cos(\omega_a \, t + \phi \cdot \phi_{BO}(t) \,) \, \cdot \, N_{\text{CBO}}(t) \cdot \, N_{\text{VW}}(t) \cdot \, N_y(t) \cdot \, N_{2\text{CBO}}(t) \cdot \, J(t) \right. \\ \left. A_{\text{BO}}(t) = 1 + A_A \cos(\omega_{\text{CBO}}(t) + \phi_A) e^{-\frac{t}{\gamma_{\text{CBO}}}} \right. \\ \left. \phi_{\text{BO}}(t) = 1 + A_\phi \cos(\omega_{\text{CBO}}(t) + \phi_\phi) e^{-\frac{t}{\gamma_{\text{CBO}}}} \right. \\ \left. N_{\text{CBO}}(t) = 1 + A_{\text{CBO}} \cos(\omega_{\text{CBO}}(t) + \phi_{\text{CBO}}) e^{-\frac{t}{\gamma_{\text{CBO}}}} \right. \\ \left. N_{\text{2CBO}}(t) = 1 + A_{2\text{CBO}} \cos(2\omega_{\text{CBO}}(t) + \phi_{2\text{CBO}}) e^{-\frac{t}{\gamma_{\text{CBO}}}} \right. \\ \left. N_{\text{VW}}(t) = 1 + A_{\text{VW}} \cos(\omega_{\text{VW}}(t) t + \phi_{\text{VW}}) e^{-\frac{t}{\gamma_{\text{VW}}}} \right. \\ \left. N_y(t) = 1 + A_y \cos(\omega_y(t) t + \phi_y) e^{-\frac{t}{\gamma_y}} \right. \\ \left. J(t) = 1 - k_{LM} \int_{t_0}^t \Lambda(t) dt \right. \\ \left. \omega_{\text{CBO}}(t) = \omega_0 t + A e^{-\frac{t}{\gamma_A}} + B e^{-\frac{t}{\gamma_B}} \right. \\ \left. \omega_y(t) = F \omega_{\text{CBO}(t)} \sqrt{2\omega_c/F} \omega_{\text{CBO}}(t) - 1 \right. \\ \left. \omega_{\text{VW}}(t) = \omega_c - 2\omega_y(t) \end{split}$$

# **Measuring the Muon Anomaly**



# **Systematics from Run-1**

$$a_{\mu} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e + C_p + C_{ml} + C_{pa}\right)}{f_{\text{calib}} \left\langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \right\rangle \left(1 + B_k + B_q\right)}$$

 $f_{\rm clock}$  •Blinded clock

 $\omega_a^m$  •Measured precession frequency

 $C_e$  •Electric field correction

 $C_p$  •Pitch correction

 $C_{ml}$  •Muon loss correction

 $C_{na}$  •Phase-acceptance correction

Absolute magnetic field calibration

 $\omega_p'(x,y,\phi)$  •Field tracking multipole distribution

 $M(x, y, \phi)$  •Muon weighted multipole distributed

 $B_k$  •Transient field from the eddy current in kicker

Bq •Transient field from the quad charging

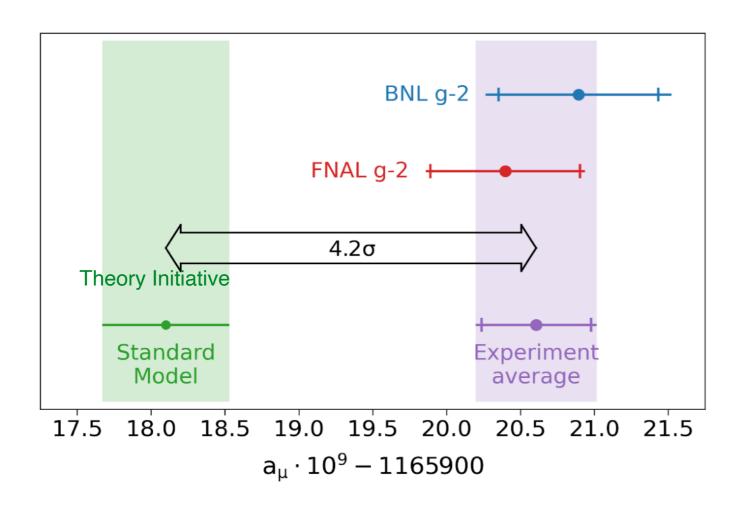
Phase acceptance and transient field corrections are the largest systematics!

 $f_{\text{calib}}$ 

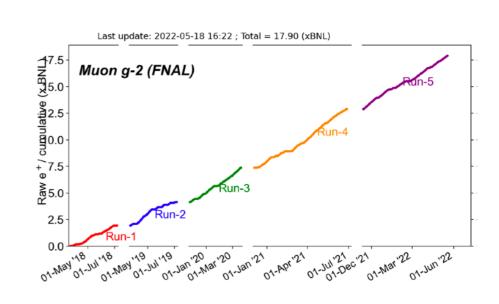
# **Crunching the numbers...**

$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	Quantity	Correction Terms	Uncertainty	
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$		(ppb)	(ppb)	Dominated by statistical
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	$\omega_a$ (statistical)	_	434	Dominated by Statistical
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	,	_	56	uncertainty
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	$C_e$	489	53	
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	$C_p$	180	13	
$\begin{array}{c ccccccccccccccccccccccccccccccccccc$	$C_{ml}$	-11	5	
$B_k$ $-27$ $37$ PA and Field $\mu'_p(34.7^\circ)/\mu_e$ $-$ 10 Transients $m_\mu/m_e$ $-$ 22 $g_e/2$ $-$ 0 Total systematic $-$ 157 *Nearly half of BNL Total fundamental factors $-$ 25 *Will be even better	$C_{pa}$	-158	75	Systematics
$B_k$ $-27$ $37$ $92$ PA and Field $\mu_p'(34.7^\circ)/\mu_e$ $-$ 10 Transients $m_\mu/m_e$ $-$ 22 $g_e/2$ $-$ 0 *Nearly half of BNL Total fundamental factors $-$ 25 *Will be even better	$f_{\rm calib}\langle\omega_p'(x,y,\phi)\times M(x,y,\phi)\rangle$	_		dominated by
$\mu_p'(34.7^\circ)/\mu_e$ — 10 Transients $m_\mu/m_e$ — 22 $g_e/2$ — 0 *Nearly half of BNL Total fundamental factors — 25 *Will be even better	$B_{m k}$	-27	37	, and the second
$m_{\mu}/m_{e}$ - 22 $g_{e}/2$ - 0  Total systematic - 157  Total fundamental factors - 25  *Will be even better	$B_{q}$	-17	92	PA and Field
$m_{\mu}/m_e$ — 22 $g_e/2$ — 0  Total systematic — 157 — *Nearly half of BNL Total fundamental factors — 25 *Will be even better	$\mu_p'(34.7^{\circ})/\mu_e$	_	10	Transients
Total systematic − 157 ← *Nearly half of BNL  Total fundamental factors − 25 *Will be even better	1	_	22	
Total systematic − 157 ← *Nearly half of BNL  Total fundamental factors − 25 *Will be even better	$g_e/2$	_	0	
"Will be even better		_	157	*Nearly half of BNL
Totals 544 462 Will be even belief	Total fundamental factors	_	25	•
100015	Totals	544	462	vviii be everi better for

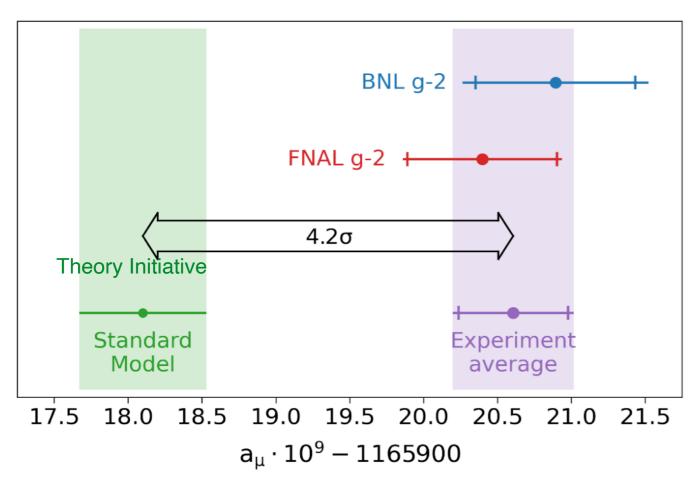
# Fermilab Muon g-2 First Result and the Comparison



- ✓ FNAL determined anomaly with 460 ppb precision
- ✓ Nothing was found that indicated contradiction with BNL results
- √15% smaller error
- √Good agreement

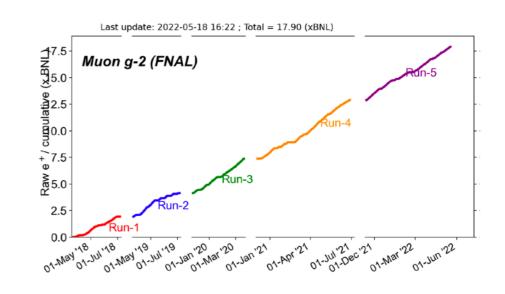


# Fermilab Muon g-2 First Result and the Comparison

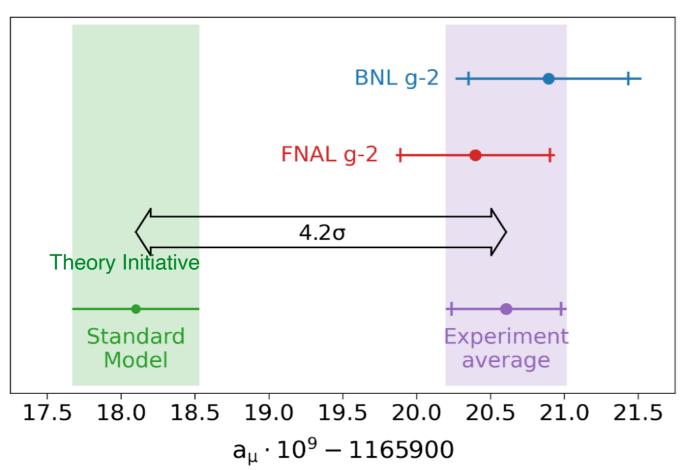


 $a_u$  (BNL)=116 592 089(63)×10<sup>-11</sup> (540 ppb)

- ✓ FNAL determined anomaly with 460 ppb precision
- ✓ Nothing was found that indicated contradiction with BNL results
- √15% smaller error
- √Good agreement



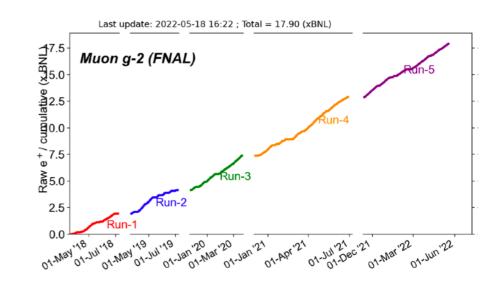
# Fermilab Muon g-2 First Result and the Comparison

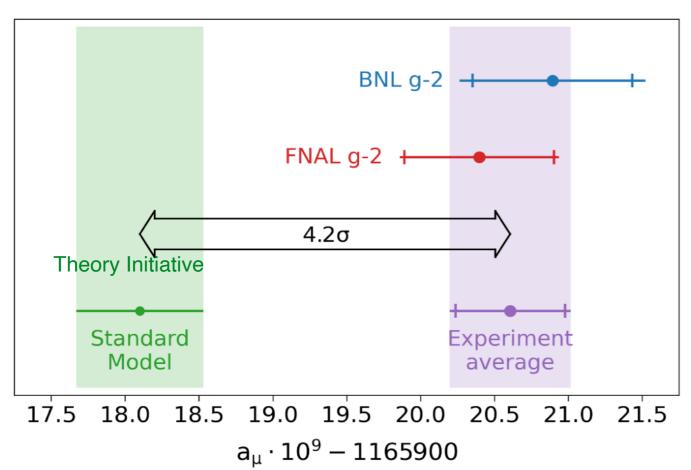


 $a_u$  (BNL)=116 592 089(63)×10<sup>-11</sup> (540 ppb)

 $a_{\mu}$  (FNAL)=116 592 040(54)×10<sup>-11</sup> (460 ppb)

- ✓ FNAL determined anomaly with 460 ppb precision
- ✓ Nothing was found that indicated contradiction with BNL results
- √15% smaller error
- √Good agreement



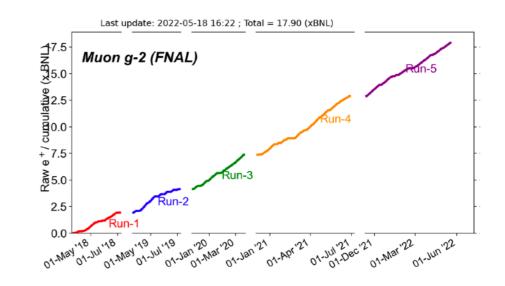


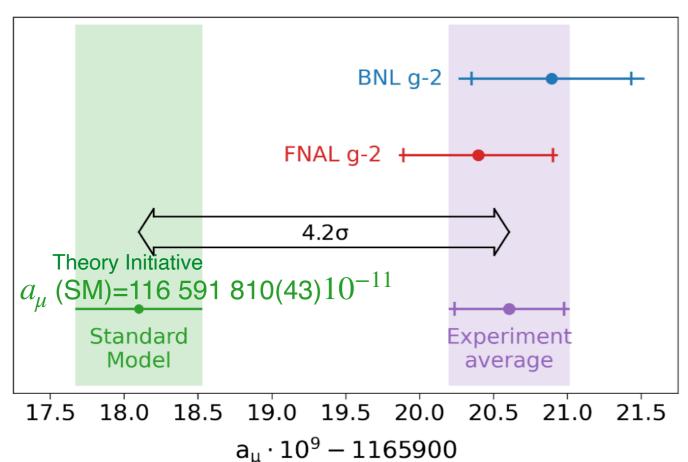
$$a_{\mu}$$
 (BNL)=116 592 089(63)× $10^{-11}$  (540 ppb)

$$a_{\mu}$$
 (FNAL)=116 592 040(54)× $10^{-11}$  (460 ppb)

$$a_{\mu}$$
 (Exp)=116 592 061(41) x $10^{-11}$ (350 ppb)

- ✓ FNAL determined anomaly with 460 ppb precision
- ✓ Nothing was found that indicated contradiction with BNL results
- √15% smaller error
- √Good agreement



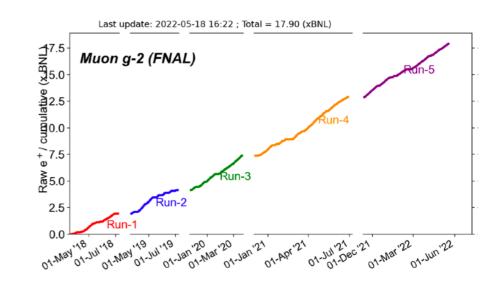


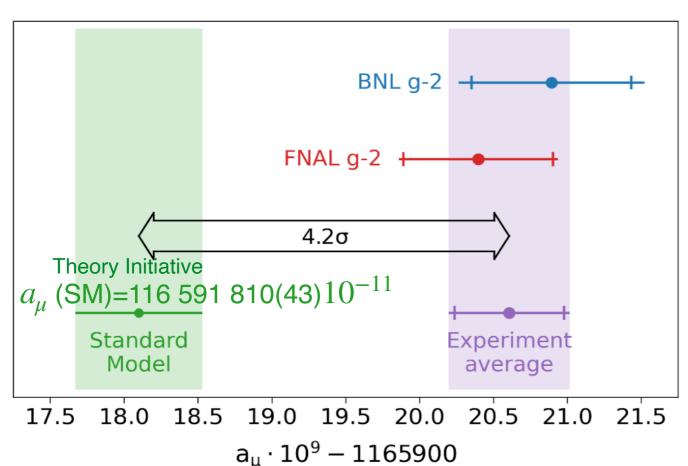
```
a_u (BNL)=116 592 089(63)×10^{-11} (540 ppb)
```

$$a_{\mu}$$
 (FNAL)=116 592 040(54)× $10^{-11}$  (460 ppb)

$$a_{\mu}$$
 (Exp)=116 592 061(41) x10<sup>-11</sup>(350 ppb)

- ✓ FNAL determined anomaly with 460 ppb precision
- ✓ Nothing was found that indicated contradiction with BNL results
- √15% smaller error
- √Good agreement





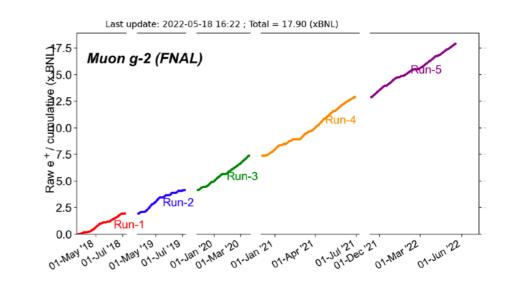
$$a_u$$
 (BNL)=116 592 089(63)× $10^{-11}$  (540 ppb)

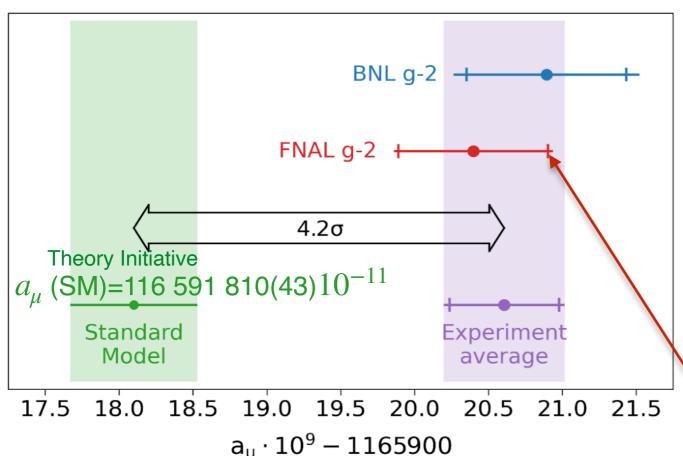
$$a_{\mu}$$
 (FNAL)=116 592 040(54)×10<sup>-11</sup> (460 ppb)

$$a_{\mu}$$
 (Exp-SM)=251(59)× $10^{-11}$  -> 4.2 $\sigma$ 

$$a_u$$
 (Exp)=116 592 061(41) x10<sup>-11</sup>(350 ppb)

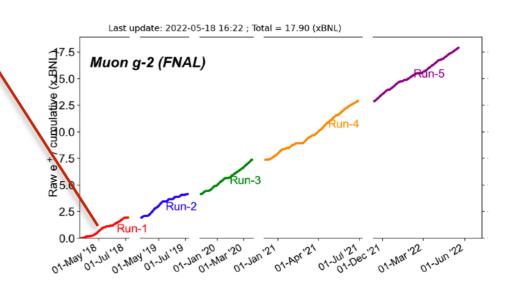
- ✓ FNAL determined anomaly with 460 ppb precision
- ✓ Nothing was found that indicated contradiction with BNL results
- √15% smaller error
- √Good agreement

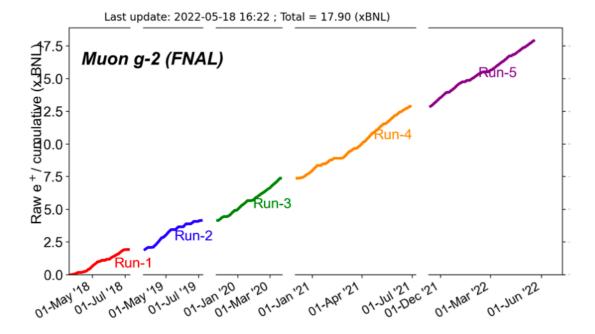


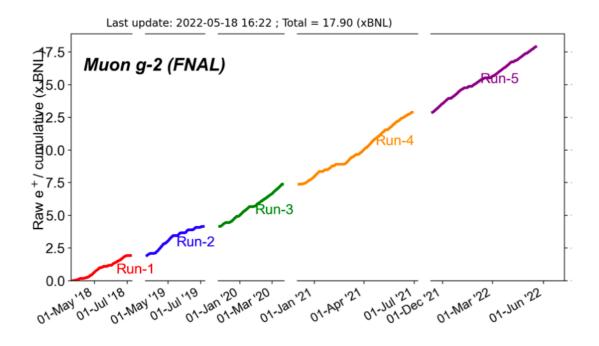


- $a_u$  (BNL)=116 592 089(63)×10<sup>-11</sup> (540 ppb)
- $a_{\mu}$  (FNAL)=116 592 040(54)×10<sup>-11</sup> (460 ppb)
- $a_u$  (Exp-SM)=251(59)× $10^{-11}$  -> 4.2 $\sigma$
- $a_{\mu}$  (Exp)=116 592 061(41) x10<sup>-11</sup>(350 ppb)

- ✓ FNAL determined anomaly with 460 ppb precision
- ✓ Nothing was found that indicated contradiction with BNL results
- √15% smaller error
- √Good agreement

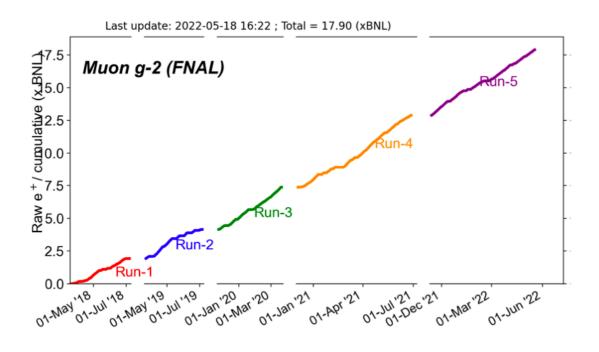






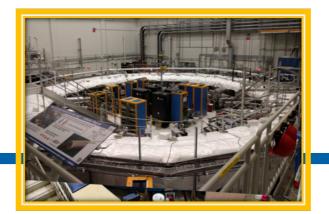
- Run-5 started in November 2021
  - Improved kick: Most recent part of Run-3 had a perfectly centered beam owing to improved kicker system.

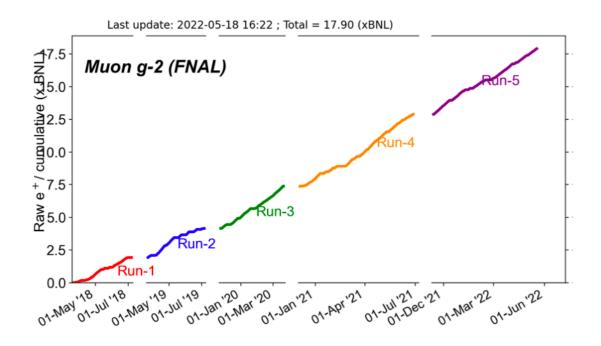




- Run-5 started in November 2021
  - Improved kick: Most recent part of Run-3 had a perfectly centered beam owing to improved kicker system.
  - Improved field stability:More stable temperature and better magnet insulation.



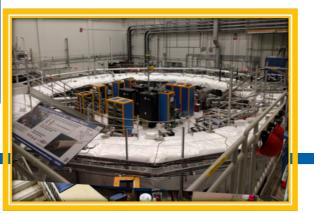


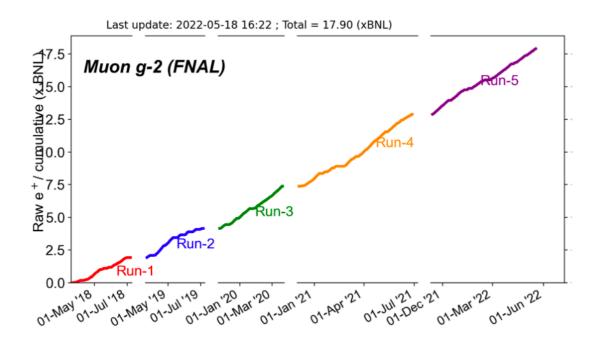






- Run-5 started in November 2021
  - Improved kick: Most recent part of Run-3 had a perfectly centered beam owing to improved kicker system.
  - Improved field stability:More stable temperature and better magnet insulation.
  - Damaged HV resistors were replaced after Run-1

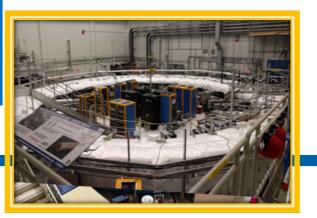


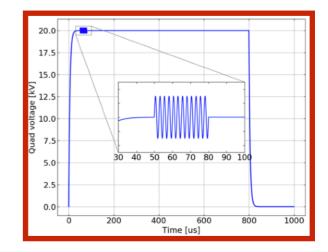






- Run-5 started in November 2021
  - Improved kick: Most recent part of Run-3 had a perfectly centered beam owing to improved kicker system.
  - Improved field stability:More stable temperature and better magnet insulation.
  - Damaged HV resistors were replaced after Run-1
  - RF System is fully integrated in Run-5

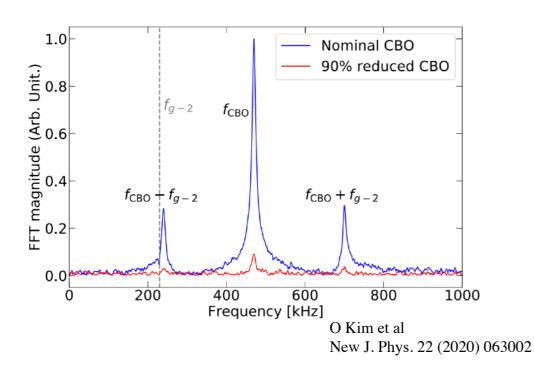


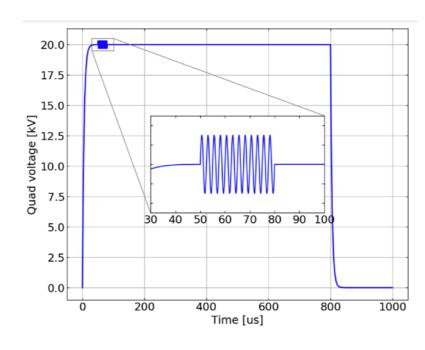


## Improvements since Run-1: Quad-RF System

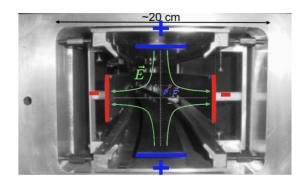
Original idea from Yuri F. Orlov and Yannis Semertzidis (Muon g-2 Note #431). By applying RF dipole or quadrupole electric field; (by kicking the beam out of phase with CBO)

- Reduce CBO Amplitude (Due to an imperfect kicker system, the beam executes CBO oscillations with a large amplitude)
- Reduce muon loss by scraping the beam



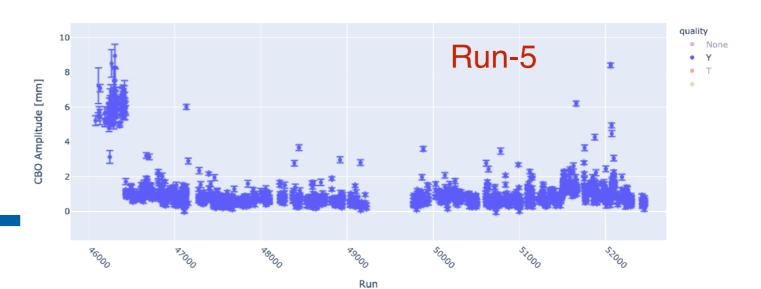


RF can be applied through the quad plates

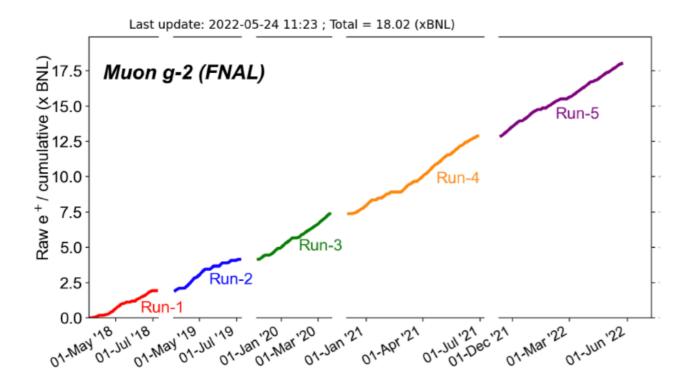


CBO Amplitude: 5-6mm->0.5-1mm

Muon Loss: factor of 4 reduction

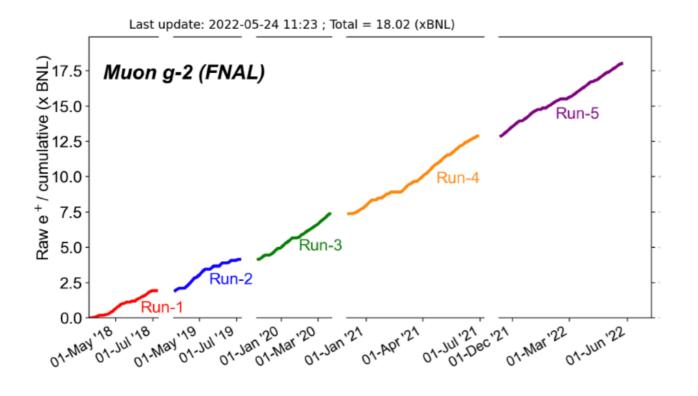


#### **Experiment Data Collection Status**



- Run-2 and Run-3 analysis are ongoing -> reduce combined exp. error by 2 times
- Run-5 will end in July-> Run-4 and beyond will reduce the statistical uncertainties down to 100 ppb in total.
- Run-6 will be negative muon running
  - Huge amount of polarity flip work will be done this summer during shutdown.

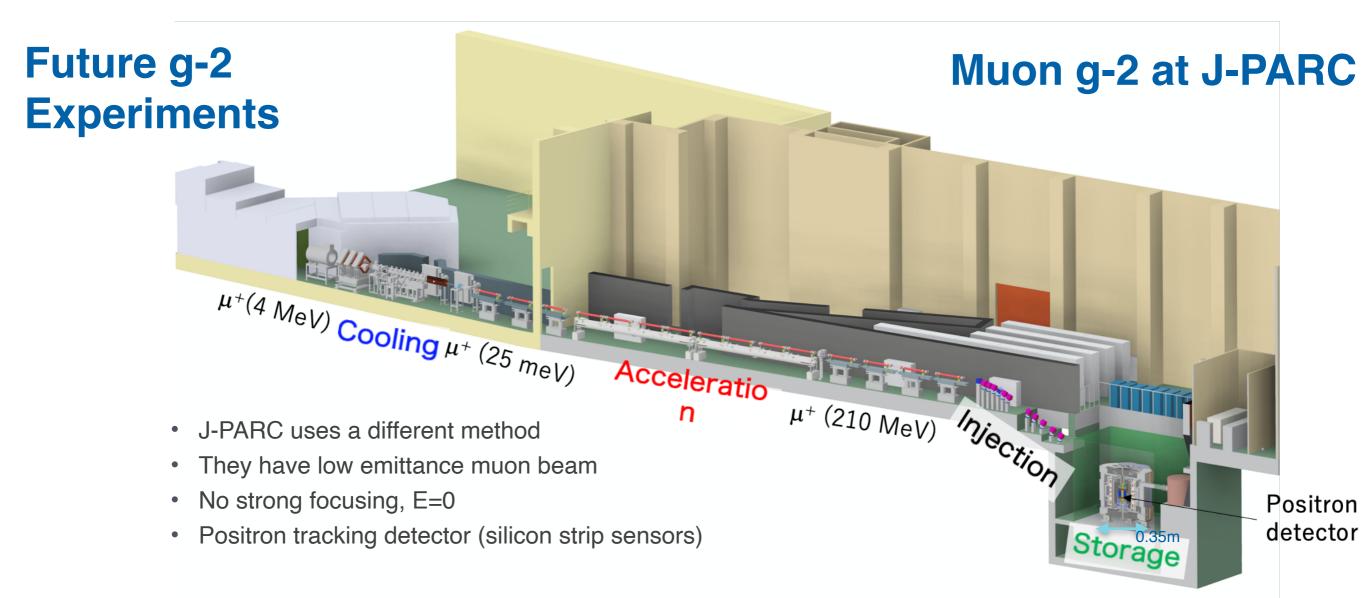
#### **Experiment Data Collection Status**



- Run-2 and Run-3 analysis are ongoing -> reduce combined exp. error by 2 times
- Run-5 will end in July-> Run-4 and beyond will reduce the statistical uncertainties down to 100 ppb in total.
- Run-6 will be negative muon running
  - Huge amount of polarity flip work will be done this summer during shutdown.

#### **Run-6 Negative Muon Run at Fermilab**

- J-PARC can only use positive muons, so probably last chance to try it in near future
- · Allow us to probe CPT and Lorentz Violation at the highest sensitivity in the muon sector
  - FNAL  $\mu^+$  & BNL  $\mu^-$
  - FNAL  $\mu^-$  & BNL  $\mu^+$
  - FNAL  $\mu^-$  & J-PARC  $\mu^+$
- Reduce the uncertainty achieved by Brookhaven in the measurement of g–2 of the negative muon by a factor of two
- Tentative plan: 7-9 months—>4xBNL



- Electric field will be eliminated by using reaccelerated thermal muon beam
- Lower momentum muon beam + compact storage region with highly uniform magnetic field
- Tracking detector for decay positrons —> reduced pile-up + able to measure the momentum direction of positrons.
- Final Goal is to reach **0.46ppm** –> **0.1ppm** on  $a_{\mu}$
- Beam line construction has started and commissioning is expected to start in 2027

## Other new ideas on measuring more precise g-2

- No future experiments other than Muon g-2 at J-PARC
- New ideas on improving the precision of magic momentum approaches -> Intense muon beam + better systematics
  - The Heavy Ion Accelerator Facility (HIAF): Intense muon source—> expected muon intensity is 30 times more than Fermilab Muon g-2 beam line—> x2 improvement on total uncertainty.
  - High-intensity muon beamline (HiMB) project at PSI: Aims to upgrade production targets and beam lines -> Proposed be used for muon EDM search to measure g-2—> Possible to reach to 0.1 ppm in 1 y
  - Non-continuous and nonuniform magnets ideas: Different B field measurement + 15GeV muon beam—> Increased measurement time —> Increased precision

#### **Summary**

- FNAL determined  $a_{\mu}$  with 460 ppb precision and confirmed BNL experimental result.
  - Run-2 and Run-3 results are expected to be published soon
  - Aiming for 100 ppb syst and 100 ppb stat. with the new improvements.
  - Getting prepared to run with  $\mu^-$  for Run-6
    - Will reduce the uncertainty achieved by Brookhaven in the measurement of g-2 of the negative muon by a factor of two
    - Improve the opportunity of probing CPT and Lorentz Violation
- J-PARC preparation is ongoing, expect to start data taking in 2027
  - \_ Goal to reduce the uncertainty on  $\,a_{\mu}$  to 0.1 ppm



# Thanks!

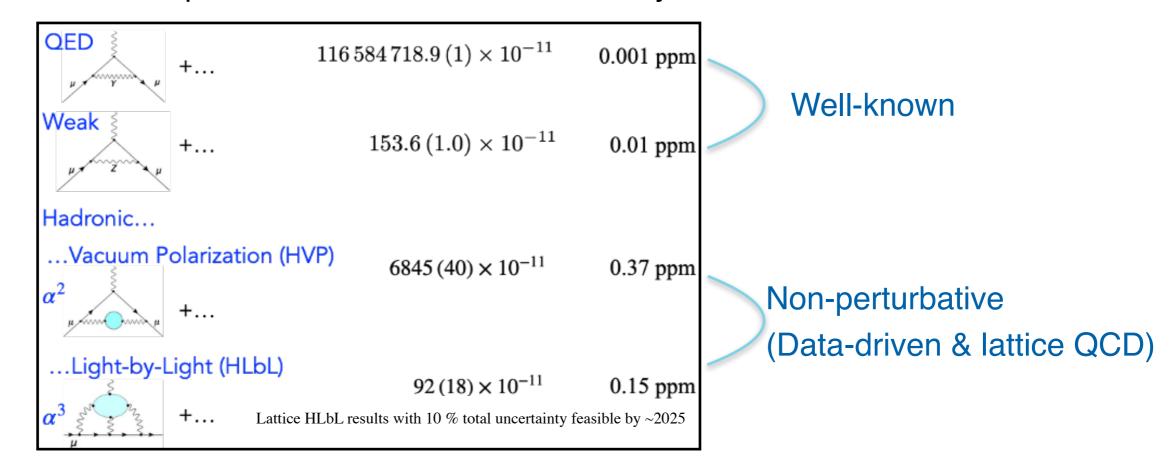
# **Back-up**

#### **Muon Magnetic Moment and Standard Model**

- Dirac predicted g-factor as 2 for spin 1/2 particles
- But then Kusch & Foley measured electron spin factor as 2.00238
- Julian Schwinder calculated this deviation in QED (1948)

$$a = (g - 2)/2 = \alpha/2\pi = 0.0016$$
, where  $\alpha = 1/137$ 

#### More loops of calculations are added today



#### **Lattice QCD Calculations**

#### Lattice

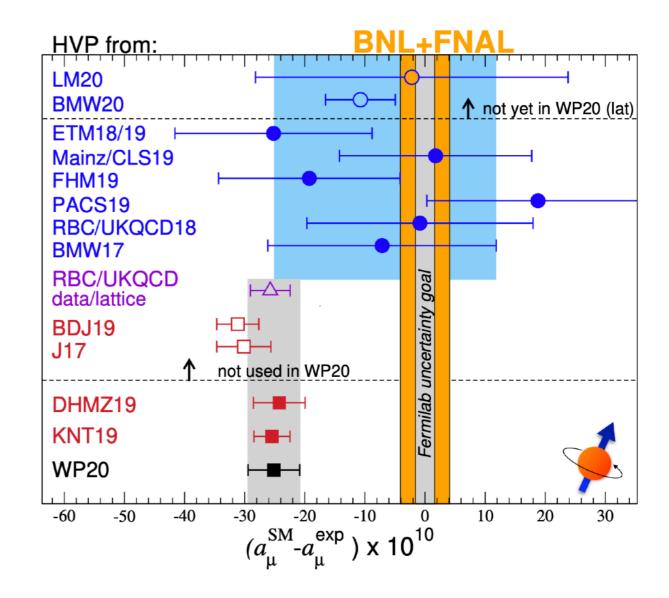
#### **Data-based Dispersive**

#### Lattice and Data

#### Official WP20

The anomalous magnetic moment of the muon in the Standard Model

T. Aoyama<sup>1,2,3</sup>, N. Asmussen<sup>4</sup>, M. Benayoun<sup>5</sup>, J. Bijnens<sup>6</sup>, T. Blum<sup>7,8</sup>, M. Bruno<sup>9</sup>, I. Caprini<sup>10</sup> C. M. Carloni Calame<sup>11</sup>, M. Cè<sup>9,12,13</sup>, G. Colangelo<sup>†14</sup>, F. Curciarello<sup>15,16</sup>, H. Czyż<sup>17</sup>, I. Danilkin<sup>12</sup>, M. Davier<sup>†18</sup> C. T. H. Davies<sup>19</sup>, M. Della Morte<sup>20</sup>, S. I. Eidelman<sup>†21,22</sup>, A. X. El-Khadra<sup>†23,24</sup>, A. Gérardin<sup>25</sup>, D. Giusti<sup>26,27</sup> M. Golterman<sup>28</sup>, Steven Gottlieb<sup>29</sup>, V. Gülpers<sup>30</sup>, F. Hagelstein<sup>14</sup>, M. Hayakawa<sup>31,2</sup>, G. Herdoíza<sup>32</sup>, D. W. Hertzog<sup>33</sup> A. Hoecker<sup>34</sup>, M. Hoferichter 14,35, B.-L. Hoid<sup>36</sup>, R. J. Hudspith 12,13, F. Ignatov<sup>21</sup>, T. Izubuchi 37,8, F. Jegerlehner L. Jin<sup>7,8</sup>, A. Keshavarzi<sup>39</sup>, T. Kinoshita<sup>40,41</sup>, B. Kubis<sup>36</sup>, A. Kupich<sup>21</sup>, A. Kupść<sup>42,43</sup>, L. Laub<sup>14</sup>, C. Lehner<sup>†26,37</sup> L. Lellouch<sup>25</sup>, I. Logashenko<sup>21</sup>, B. Malaescu<sup>5</sup>, K. Maltman<sup>44,45</sup>, M. K. Marinković<sup>46,47</sup>, P. Masjuan<sup>48,49</sup> A. S. Meyer<sup>37</sup>, H. B. Meyer<sup>12,13</sup>, T. Mibe<sup>†1</sup>, K. Miura<sup>12,13,3</sup>, S. E. Müller<sup>50</sup>, M. Nio<sup>2,51</sup>, D. Nomura<sup>52,53</sup>, A. Nyffeler 12, V. Pascalutsa 2, M. Passera 4, E. Perez del Rio 5, S. Peris 48,49, A. Portelli 0, M. Procura 6 C. F. Redmer<sup>12</sup>, B. L. Roberts<sup>†57</sup>, P. Sánchez-Puertas<sup>49</sup>, S. Serednyakov<sup>21</sup>, B. Shwartz<sup>21</sup>, S. Simula<sup>27</sup> D. Stöckinger<sup>58</sup>, H. Stöckinger-Kim<sup>58</sup>, P. Stoffer<sup>59</sup>, T. Teubner<sup>‡60</sup>, R. Van de Water<sup>24</sup>, M. Vanderhaeghen<sup>12,13</sup> G. Venanzoni<sup>61</sup>, G. von Hippel<sup>12</sup>, H. Wittig<sup>12,13</sup>, Z. Zhang<sup>18</sup> M. N. Achasov<sup>21</sup>, A. Bashir<sup>62</sup>, N. Cardoso<sup>47</sup>, B. Chakraborty<sup>63</sup>, E.-H. Chao<sup>12</sup>, J. Charles<sup>25</sup>, A. Crivellin<sup>64,6</sup> O. Deineka<sup>12</sup>, A. Denig<sup>12,13</sup>, C. DeTar<sup>66</sup>, C. A. Dominguez<sup>67</sup>, A. E. Dorokhov<sup>68</sup>, V. P. Druzhinin<sup>21</sup>, G. Eichmann<sup>69,47</sup> M. Fael<sup>70</sup>, C. S. Fischer<sup>71</sup>, E. Gámiz<sup>72</sup>, Z. Gelzer<sup>23</sup>, J. R. Green<sup>9</sup>, S. Guellati-Khelifa<sup>73</sup>, D. Hatton<sup>19</sup>, N. Hermansson-Truedsson<sup>14</sup>, S. Holz<sup>36</sup>, B. Hörz<sup>74</sup>, M. Knecht<sup>25</sup>, J. Koponen<sup>1</sup>, A. S. Kronfeld<sup>24</sup>, J. Laiho<sup>75</sup> S. Leupold<sup>42</sup>, P. B. Mackenzie<sup>24</sup>, W. J. Marciano<sup>37</sup>, C. McNeile<sup>76</sup>, D. Mohler<sup>12,13</sup>, J. Monnard<sup>14</sup>, E. T. Neil<sup>77</sup> A. V. Nesterenko<sup>68</sup>, K. Ottnad<sup>12</sup>, V. Pauk<sup>12</sup>, A. E. Radzhabov<sup>78</sup>, E. de Rafael<sup>25</sup>, K. Raya<sup>79</sup>, A. Risch<sup>12</sup>. A. Rodríguez-Sánchez<sup>6</sup>, P. Roig<sup>80</sup>, T. San José<sup>12,13</sup>, E. P. Solodov<sup>21</sup>, R. Sugar<sup>81</sup>, K. Yu. Todyshev<sup>21</sup>, A. Vainshtein<sup>82</sup>, A. Vaquero Avilés-Casco<sup>66</sup>, E. Weil<sup>71</sup>, J. Wilhelm<sup>12</sup>, R. Williams<sup>71</sup>, A. S. Zhevlakov<sup>78</sup>



- BMW20 is also in tension with data-based dispersive result(2.1 $\sigma$ )
- Expecting results from RBC/UKQCD and FNAL/MILC in the coming months
- Will be interested to see how it evolves in future but lattice QCD calculations requires a
  huge amount of computing resource. But all groups are working on defining intermediate
  results (simpler way to compare things)
- Lattice HVP (average) with ≤ 0.5 % errors feasible by 2025

RPF Spring meeting by A. El-Khadra

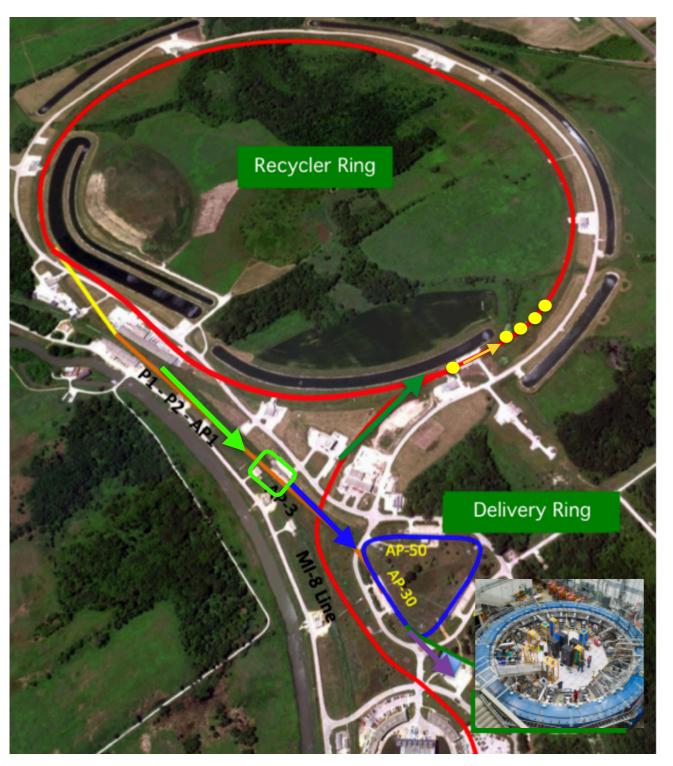
#### Most promising models to explain the discrepancy

- Muon g-2 can indicate if there is a CP-conserving, lepton-flour conserving or BSM chirality-flipping interaction but can't tell which one is the most promising.
- Possible explanations:
  - SUSY models (while evading LHC limits)
  - Leptoquark models (if leptoquark masses are above all LHC limits)
  - 2-Higgs doublet models



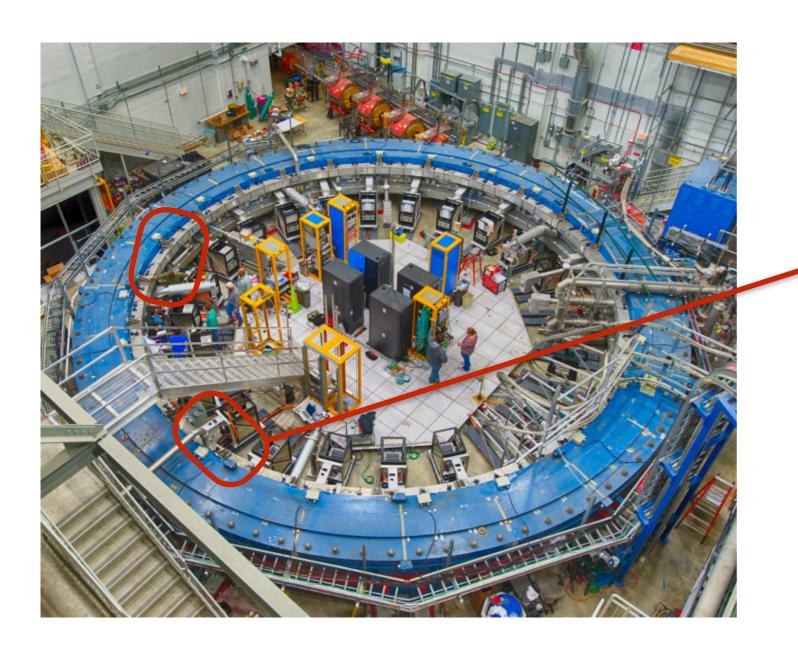
Establishing a g-2 discrepancy from SM would place a strict limit on BSM scenarios

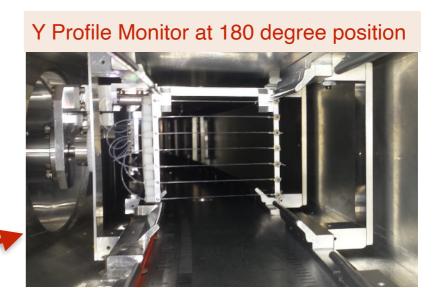
#### **Muon Campus at Fermilab**



- 8 GeV protons are delivered to Recycler Ring from Booster
- Split the proton bunch into four bunches with RF system.
- Direct the proton punches to pion production target and obtain pions.
- Muons produced by pion decays circulate in the delivery ring until proton contamination is removed.
- Deliver muons to g-2 storage ring.

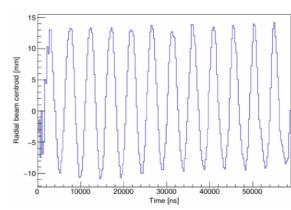
## **Monitoring the Injected Beam**



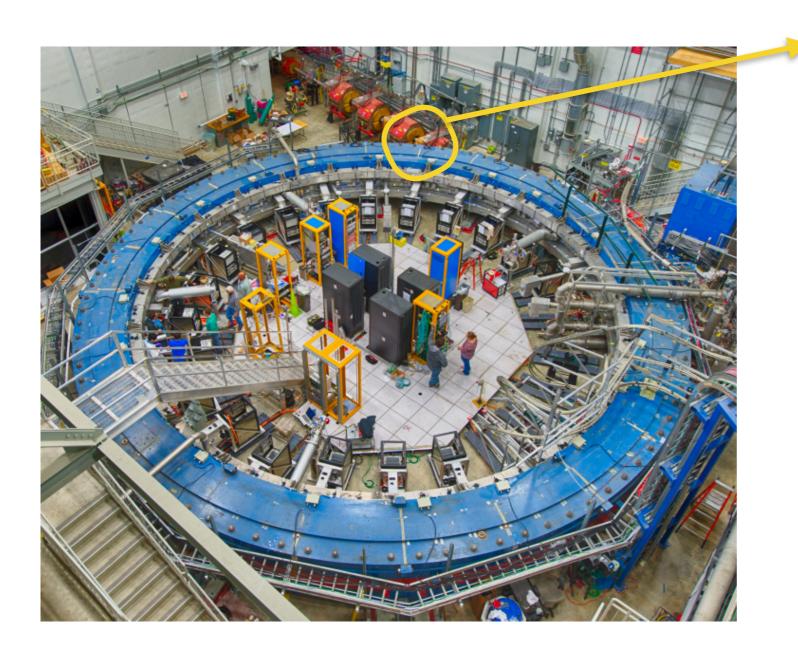


#### Fiber Harps

- Scintillating Fibers
- x and y profile monitoring
- Tool for beam commissioning

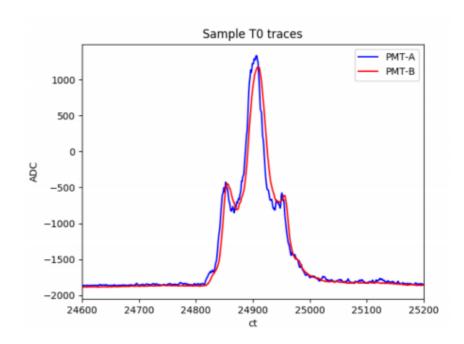


## Monitoring the Injected Beam: T0 counter



#### T0 Counter

- Thin scintillator with 2 PMT readouts
- Provides beam time profile before beam enters the storage ring

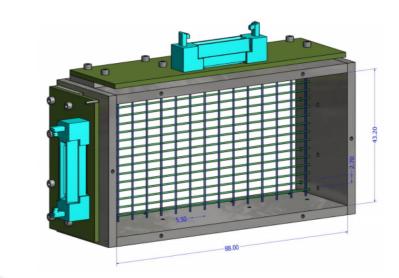


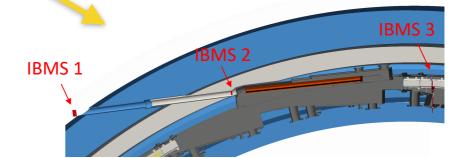
#### Monitoring the Injected Beam: IBMS detectors

IBMS 1 Y

IBMS spatial beam profiles IBMS 1 X
mean: 9.9 fibers, 13.3 mm
RMS: 2.9 fibers, 15.8 mm

- IBMS (Inflector beam monitoring systems)
  - Check beam injection characteristics
  - 2 planes of scintillation fibers





## **Quad System Upgrades**

## Refurbished and upgraded the BNL EQS



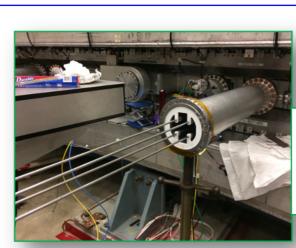




Adding HV Feedthrough extension to avoid spark in the HV components 

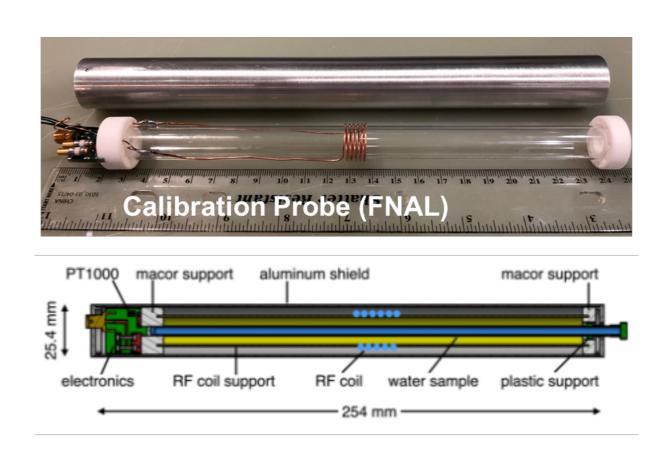
Avoid damaging the system components

New HV Pulsers were designed to pulse the EQS → Reduce the spark rate and gain more reliability on operating



Installing additional (Macor) support to fix the HV extension lead-> Prevent mechanical vibration and provide electrical isolation

# NMR Probes $\frac{a_{\mu}}{\text{Calib }} \propto \underbrace{\frac{f_{\text{clock}} \ \omega_a^m \, (1 + C_e + C_p + C_{ml} + C_{pa})}{f_{\text{calib}} \ \langle \omega_p'(x,y,\phi) \times M(x,y,\phi) \rangle (1 + B_k + B_q)}$



#### Absolute field calibration:

- Absolute probes were used to calibrate NMR probes
- Proton NMR, calibrated in terms of  $\omega_p(T_r)$  of a proton shielded in a spherical sample of water at an exact temperature.

$$\frac{a_{\mu}}{f_{\text{calib}}} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e + C_p + C_{ml} + \frac{C_{pa}}{D_q}\right)}{f_{\text{calib}} \left\langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \right\rangle \left(1 + \frac{C_{pa}}{B_k} + \frac{B_q}{D_q}\right)}$$

$$\frac{a_{\mu}}{f_{\text{calib}}} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e + C_p + C_{ml} + C_{pa}\right)}{f_{\text{calib}} \left\langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \right\rangle \left(1 + B_k + B_q\right)}$$



$$\frac{a_{\mu}}{f_{\text{calib}}} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e + C_p + C_{ml} + C_{pa}\right)}{f_{\text{calib}} \left\langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \right\rangle \left(1 + B_k + B_q\right)}$$





$$\frac{a_{\mu}}{f_{\text{calib}}} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e + C_p + C_{ml} + \frac{C_{pa}}{D_q}\right)}{f_{\text{calib}} \left\langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \right\rangle \left(1 + \frac{C_{pa}}{B_k}\right)}$$





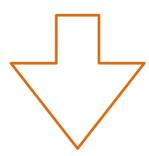
Field Perturbation

$$\frac{a_{\mu}}{f_{\text{calib}}} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e + C_p + C_{ml} + \frac{C_{pa}}{C_{pa}}\right)}{f_{\text{calib}} \left\langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \right\rangle \left(1 + \frac{B_k + B_q}{B_q}\right)}$$





Field Perturbation



$$a_{\mu} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e + C_p + C_{ml} + \frac{C_{pa}}{C_{pa}}\right)}{f_{\text{calib}} \ \langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \rangle \ \left(1 + \frac{B_k + B_q}{B_q}\right)}$$

#### Kickers pulsing created influence on the average field seen by beam

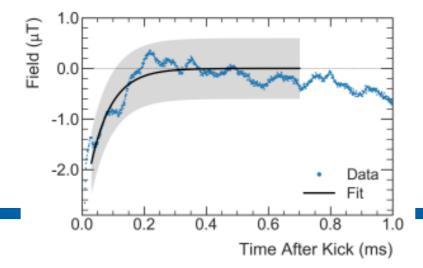




Field Perturbation

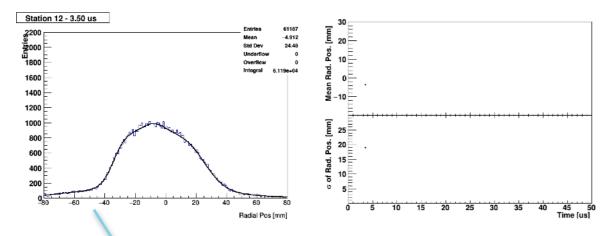


#### Used a magnetometer to measure the transient

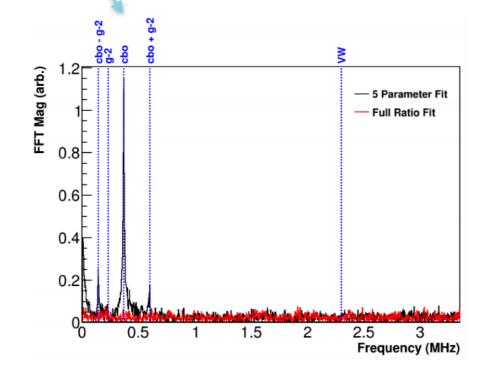


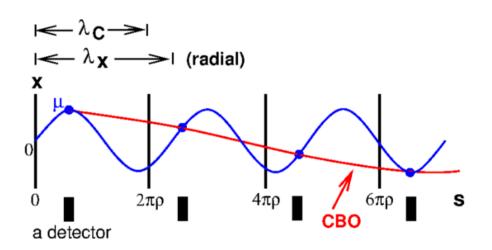
#### **Beam Dynamics**

CBO - Coherent Beam Oscillation



Radial CBO movement





 $\lambda_x$  - radial wavelength

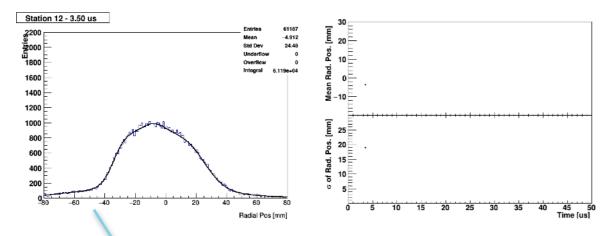
 $\lambda_C$  - cyclotron wavelength

Frequency from detector point of view =  $f_c - f_x$ 

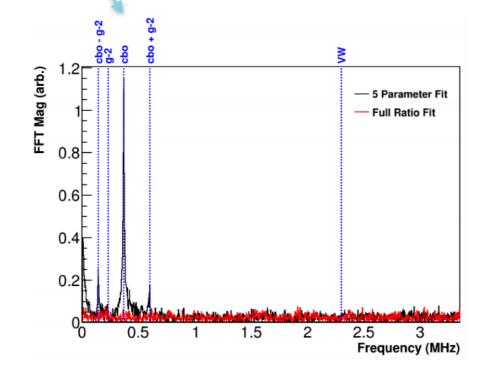
$$x=x_e+A_xcos(f_xt+\delta_x)$$
  $y=A_ycos(f_yt+\delta_y)$  Simple Harmonic Motion

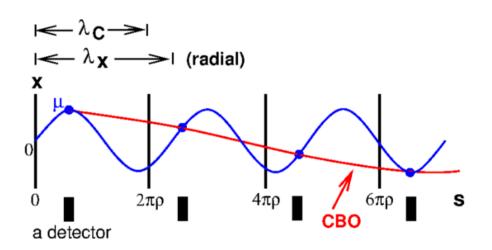
#### **Beam Dynamics**

CBO - Coherent Beam Oscillation



Radial CBO movement





 $\lambda_x$  - radial wavelength

 $\lambda_C$  - cyclotron wavelength

Frequency from detector point of view =  $f_c - f_x$ 

$$x=x_e+A_xcos(f_xt+\delta_x)$$
  $y=A_ycos(f_yt+\delta_y)$  Simple Harmonic Motion

## Phase Acceptance

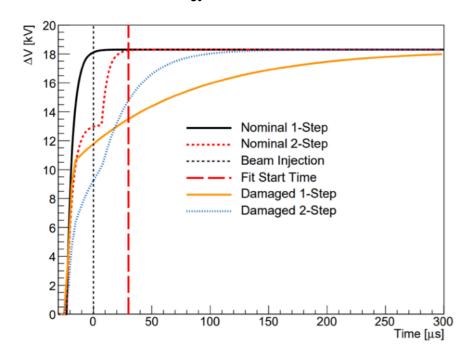
$$\frac{a_{\mu}}{f_{\text{calib}}} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e + C_p + C_{ml} + \frac{C_{pa}}{C_{pa}}\right)}{f_{\text{calib}} \left\langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \right\rangle \left(1 + B_k + B_q\right)}$$

When there is a time dependent phase, It shifts the  $\omega_a$  !

## **Phase Acceptance**

$$a_{\mu} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e + C_p + C_{ml} + C_{pa}\right)}{f_{\text{calib}} \ \langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \rangle \ \left(1 + B_k + B_q\right)}$$

When there is a time dependent phase, It shifts the  $\omega_a$  !

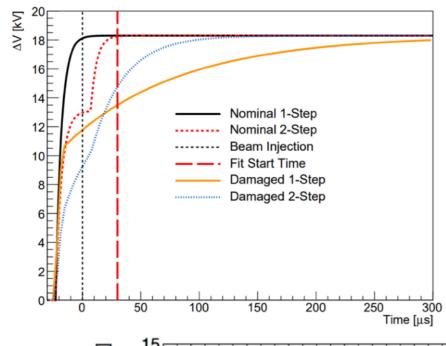


Due to damaged HV resistors; stored beam distribution was unstable.

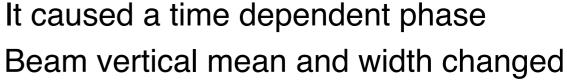
### **Phase Acceptance**

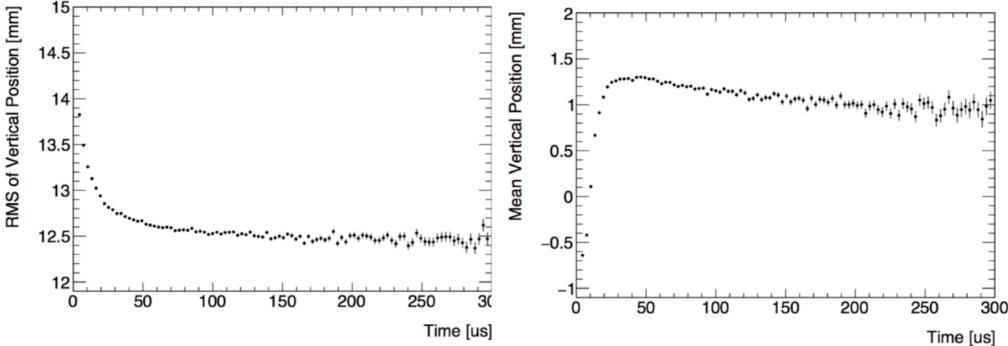
$$a_{\mu} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e + C_p + C_{ml} + C_{pa}\right)}{f_{\text{calib}} \ \langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \rangle \ \left(1 + B_k + B_q\right)}$$

When there is a time dependent phase, It shifts the  $\omega_a$  !



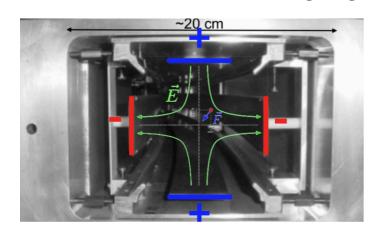
Due to damaged HV resistors; stored beam distribution was unstable.



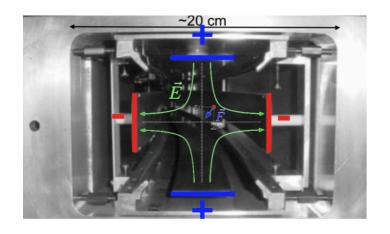


$$\frac{a_{\mu}}{f_{\text{calib}}} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e + C_p + C_{ml} + C_{pa}\right)}{f_{\text{calib}} \left\langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \right\rangle \left(1 + B_k + B_q\right)}$$

$$\frac{a_{\mu}}{f_{\text{calib}}} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e + C_p + C_{ml} + C_{pa}\right)}{f_{\text{calib}} \left\langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \right\rangle \left(1 + B_k + B_q\right)}$$

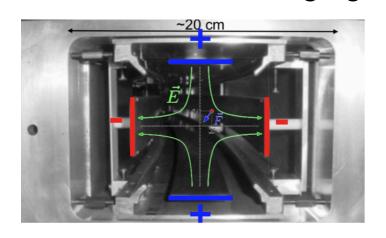


$$\frac{a_{\mu}}{f_{\text{calib}}} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e + C_p + C_{ml} + C_{pa}\right)}{f_{\text{calib}} \left\langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \right\rangle \left(1 + B_k + B_q\right)}$$



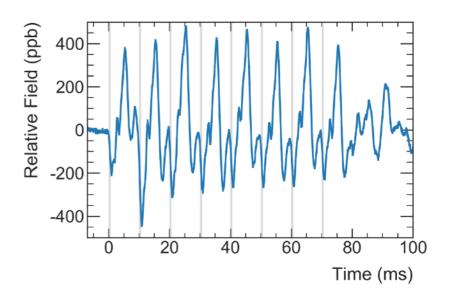


$$\frac{a_{\mu}}{f_{\text{calib}}} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e + C_p + C_{ml} + \frac{C_{pa}}{F_{\text{calib}}}\right)}{f_{\text{calib}} \left\langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \right\rangle \left(1 + B_k + \frac{B_q}{F_{q}}\right)}$$

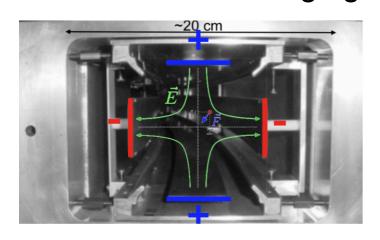




Field Perturbation



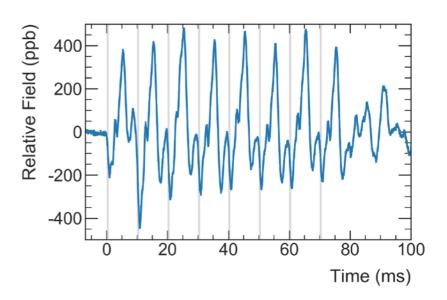
$$\frac{a_{\mu}}{f_{\text{calib}}} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e + C_p + C_{ml} + \frac{C_{pa}}{F_{\text{calib}}}\right)}{f_{\text{calib}} \left\langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \right\rangle \left(1 + B_k + \frac{B_q}{F_{q}}\right)}$$





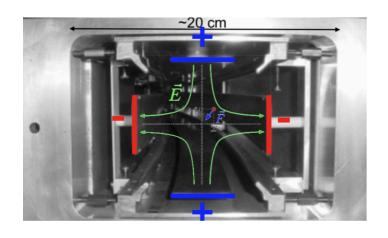
Field Perturbation





$$\frac{a_{\mu}}{f_{\text{calib}}} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e + C_p + C_{ml} + \frac{C_{pa}}{F_{\text{calib}}}\right)}{f_{\text{calib}} \left\langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \right\rangle \left(1 + B_k + \frac{B_q}{F_{q}}\right)}$$

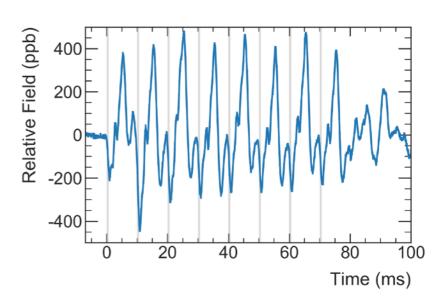
#### Quads charging and discharging cause mechanical vibration



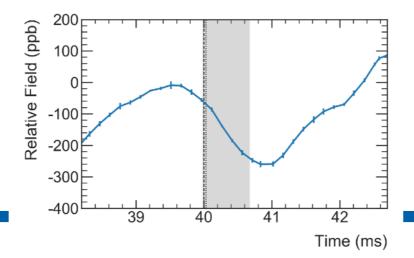


Field Perturbation





Measure the field with special NMR probes and map the effect!



$$\frac{a_{\mu}}{f_{\text{calib}}} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e \right) \left(C_p + C_{ml} + C_{pa}\right)}{f_{\text{calib}} \left\langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \right\rangle \left(1 + B_k + B_q\right)}$$

$$\vec{\omega}_{a} = \frac{e}{m} \left[ a_{\mu} \vec{B} - a_{\mu} \frac{\gamma}{\gamma + 1} (\vec{\beta} \cdot \vec{B}) \vec{\beta} - \left( a_{\mu} - \frac{1}{\gamma^{2} - 1} \right) \vec{\beta} \times \vec{E} \right]$$

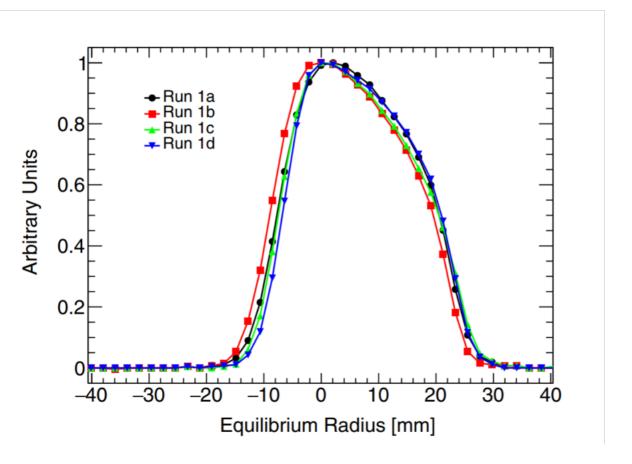
$$\frac{a_{\mu}}{f_{\text{calib}}} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e\right) \left(C_p\right) + C_{ml} + C_{pa}}{f_{\text{calib}} \ \langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \rangle \ \left(1 + B_k + B_q\right)}$$

Not all of the muons are at magic momentum! There is a 0.5% momentum acceptance

$$\vec{\omega}_{a} = \frac{e}{m} \left[ a_{\mu} \vec{B} - a_{\mu} \frac{\gamma}{\gamma + 1} (\vec{\beta} \cdot \vec{B}) \vec{\beta} - \left( a_{\mu} - \frac{1}{\gamma^{2} - 1} \right) \vec{\beta} \times \vec{E} \right]$$

$$\frac{a_{\mu}}{f_{\text{calib}}} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e\right) \left(C_p\right) + C_{ml} + C_{pa}}{f_{\text{calib}} \ \langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \rangle \ \left(1 + B_k + B_q\right)}$$

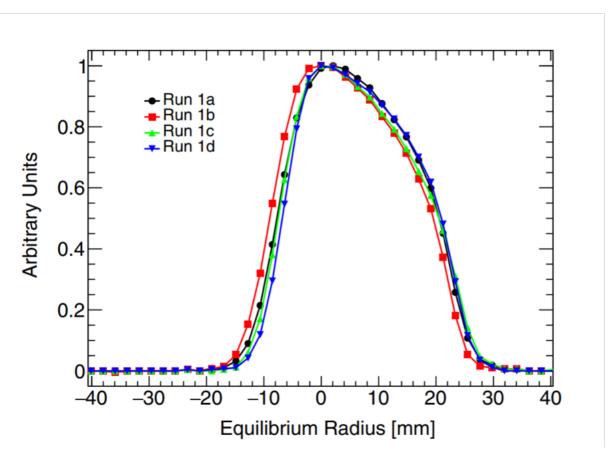
Not all of the muons are at magic momentum! There is a 0.5% momentum acceptance



$$\vec{\omega}_{a} = \frac{e}{m} \left[ a_{\mu} \vec{B} - a_{\mu} \frac{\gamma}{\gamma + 1} (\vec{\beta} \cdot \vec{B}) \vec{\beta} - \left( a_{\mu} - \frac{1}{\gamma^{2} - 1} \right) \vec{\beta} \times \vec{E} \right]$$

$$\frac{a_{\mu}}{f_{\text{calib}}} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e\right) \left(C_p\right) + C_{ml} + C_{pa}}{f_{\text{calib}} \ \langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \rangle \ \left(1 + B_k + B_q\right)}$$

Not all of the muons are at magic momentum! There is a 0.5% momentum acceptance

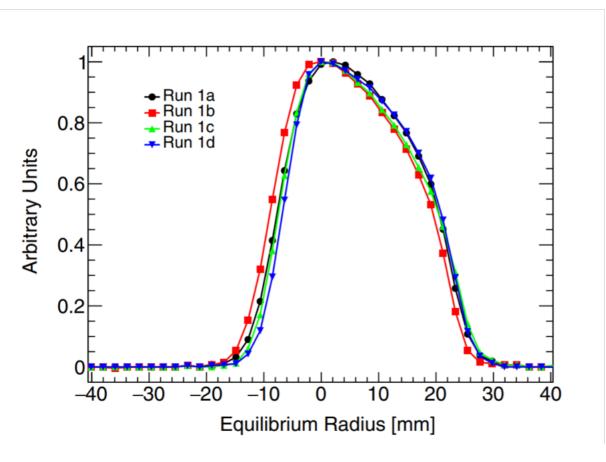


$$\vec{\omega}_{a} = \frac{e}{m} \left[ a_{\mu} \vec{B} - a_{\mu} \frac{\gamma}{\gamma + 1} (\vec{\beta} \cdot \vec{B}) \vec{\beta} - \left( a_{\mu} - \frac{1}{\gamma^{2} - 1} \right) \vec{\beta} \times \vec{E} \right]$$

Vertical betatron oscillations cause non-zero average value for  $\overrightarrow{\beta}$  .  $\overrightarrow{B}$ 

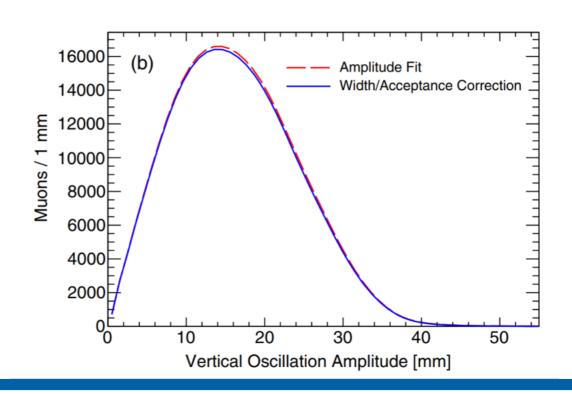
$$\frac{a_{\mu}}{f_{\text{calib}}} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e \right) \left(C_p + C_{ml} + C_{pa}\right)}{f_{\text{calib}} \left\langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \right\rangle \left(1 + B_k + B_q\right)}$$

Not all of the muons are at magic momentum! There is a 0.5% momentum acceptance



$$\vec{\omega}_{a} = \frac{e}{m} \left[ a_{\mu} \vec{B} - a_{\mu} \frac{\gamma}{\gamma + 1} (\vec{\beta} \cdot \vec{B}) \vec{\beta} - \left( a_{\mu} - \frac{1}{\gamma^{2} - 1} \right) \vec{\beta} \times \vec{E} \right]$$

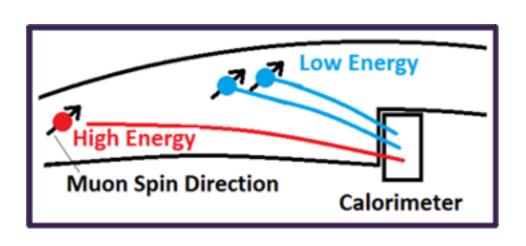
Vertical betatron oscillations cause non-zero average value for  $\overrightarrow{\beta}$  .  $\overrightarrow{B}$ 

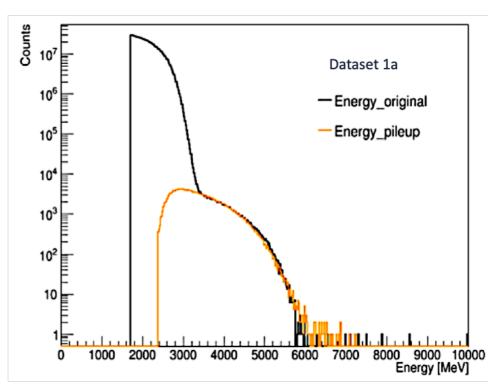


### **Pileup**

$$\frac{a_{\mu}}{f_{\text{calib}}} \propto \frac{f_{\text{clock}}(\omega_a^m) (1 + C_e + C_p + C_{ml} + C_{pa})}{f_{\text{calib}} \langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \rangle (1 + B_k + B_q)}$$

- Pileup is one of the systematics that modulated precession frequency.
- When more than two positrons hit the detector at the same time and place, they could be treated as a single pulse.
- That distorts the time and energy spectrum!





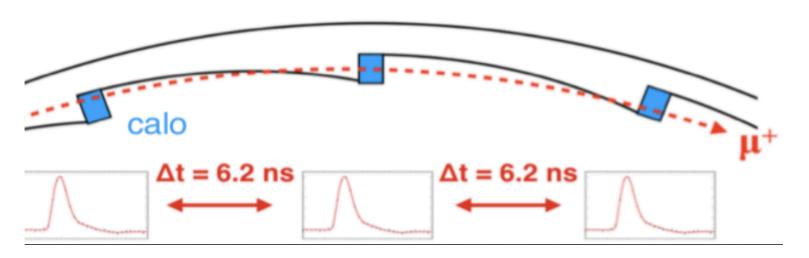
Run-1 data set	1a	1b	1c	1d
Gain changes (ppb)	12	9	9	5
Pileup (ppb)	39	42	35	31
CBO (ppb)	42	49	32	35
Time randomization (ppb)	15	12	9	7
Early-to-late effect (ppb)	21	21	22	10
total systematic uncertainty (ppb)	64	70	54	49

#### **Muon Loss**

$$\frac{a_{\mu}}{f_{\text{calib}}} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e + C_p + C_{ml} + C_{pa}\right)}{f_{\text{calib}} \left\langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \right\rangle \left(1 + B_k + B_q\right)}$$

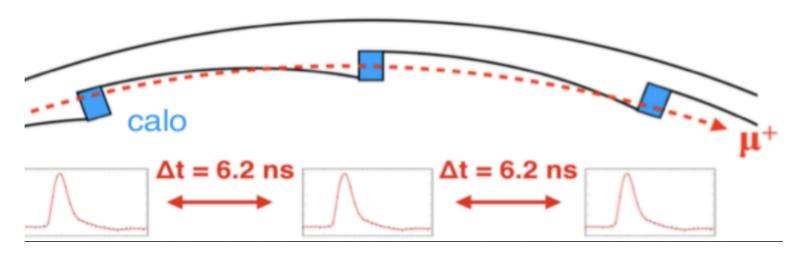
#### **Muon Loss**

$$\frac{a_{\mu}}{f_{\text{calib}}} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e + C_p + C_{ml} + C_{pa}\right)}{f_{\text{calib}} \left\langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \right\rangle \left(1 + B_k + B_q\right)}$$



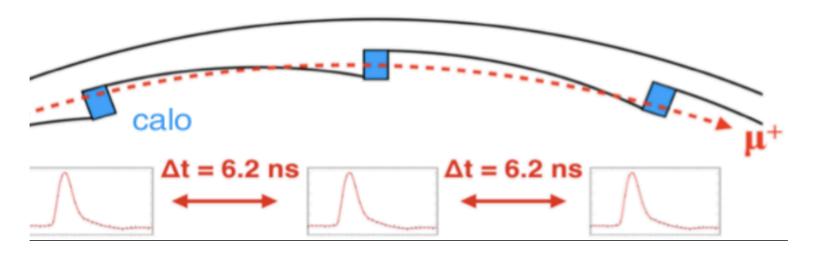
#### **Muon Loss**

$$\frac{a_{\mu}}{f_{\text{calib}}} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e + C_p + C_{ml}\right) + C_{pa}}{f_{\text{calib}} \ \langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \rangle \ (1 + B_k + B_q)}$$

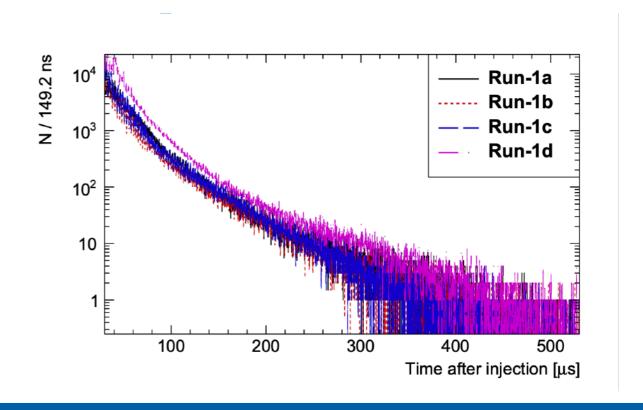


- Muons that were scattered from different materials before decaying and then punch through multiple calorimeters.
- They have different phase than stored muons so they modulate  $\omega_a$ , producing a systematic error.
- We need to identify them in the data!

$$\frac{a_{\mu}}{f_{\text{calib}}} \propto \frac{f_{\text{clock}} \ \omega_a^m \left(1 + C_e + C_p + C_{ml}\right) + C_{pa}}{f_{\text{calib}} \ \langle \omega_p'(x, y, \phi) \times M(x, y, \phi) \rangle \ \left(1 + B_k + B_q\right)}$$



- Muons that were scattered from different materials before decaying and then punch through multiple calorimeters.
- They have different phase than stored muons so they modulate  $\omega_a$ , producing a systematic error.
- We need to identify them in the data!



#### Calculating $\omega_a$ : Different analysis techniques to reduce systematic errors

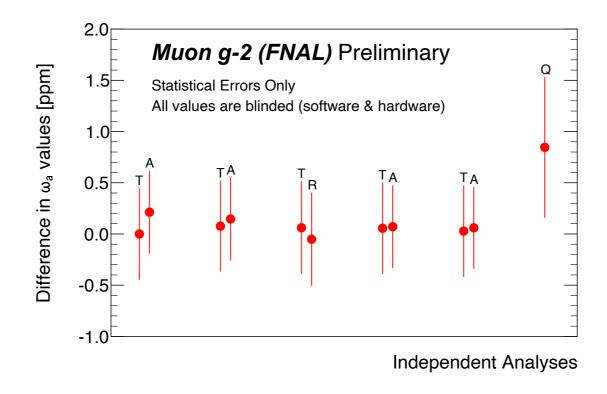


#### Fitting Methods

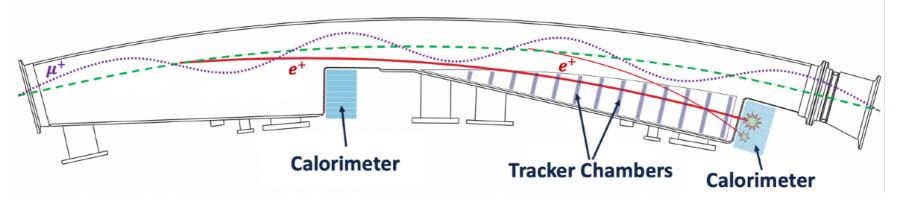
- T method:
  - Apply an energy threshold to choose
- Asymmetry-Weighted Method
  - Apply a smaller energy threshold and weight energy bin by g-2 Asymmetry.
- Ratio method:
  - Remove the exponential decay of the positron spectrum by sorting the positron time spectra into four equal subsets and then combining them again.

Requires reconstructing the positron time and energy information

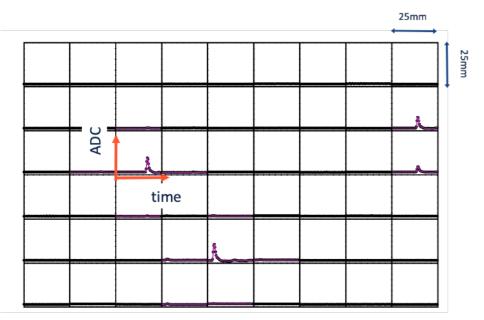
- Q-method:
  - A unique energy integrated method where you use sum over digitizer traces above noise level.

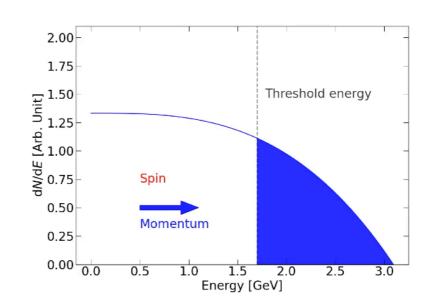


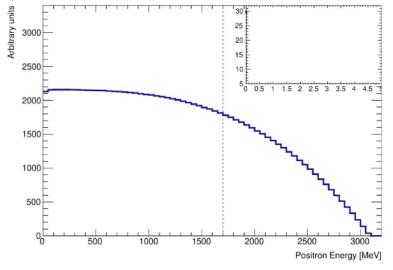




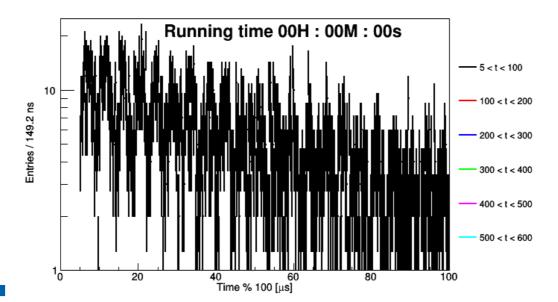
Energy distribution recorded in the calorimeter.



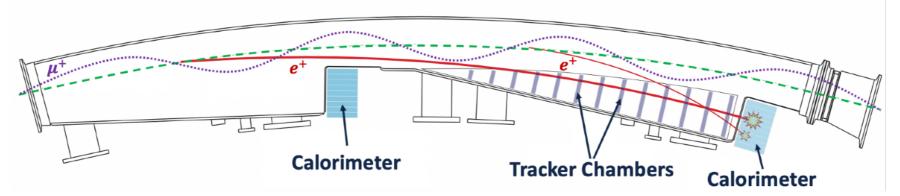




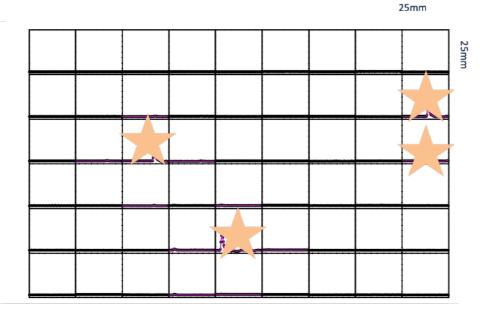
 $\omega_a$  blinded by both hardware and software

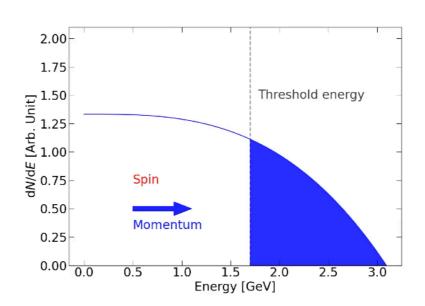


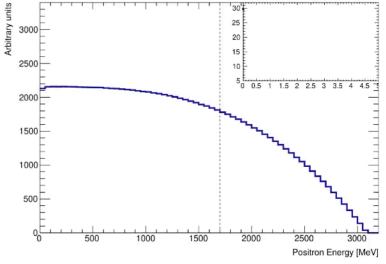




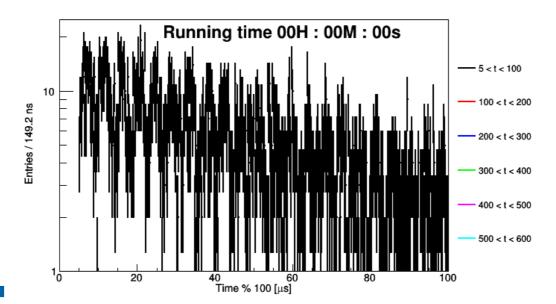
Energy distribution recorded in the calorimeter.



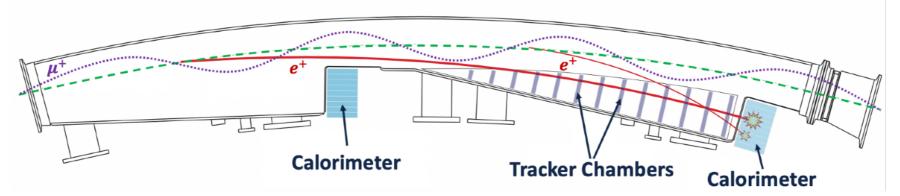




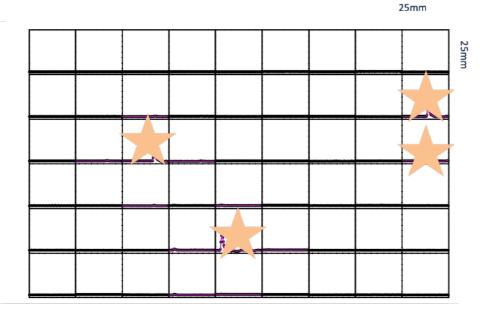
 $\omega_a$  blinded by both hardware and software

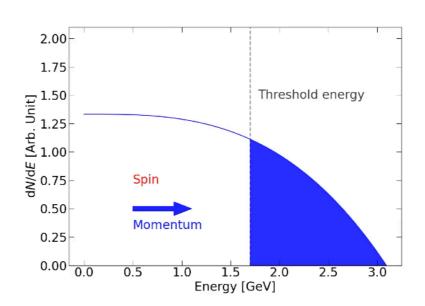


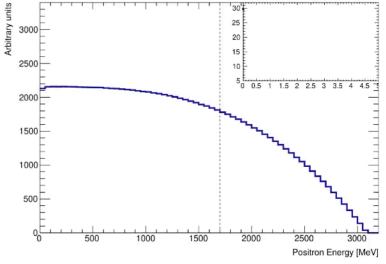




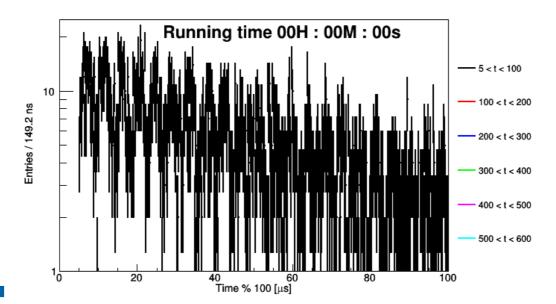
Energy distribution recorded in the calorimeter.



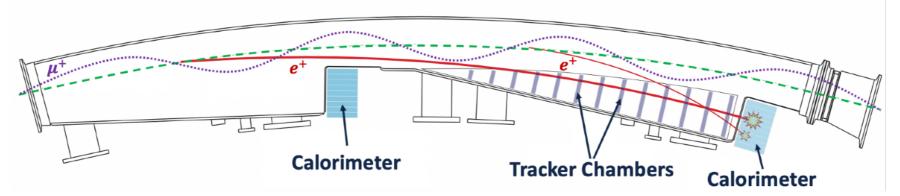




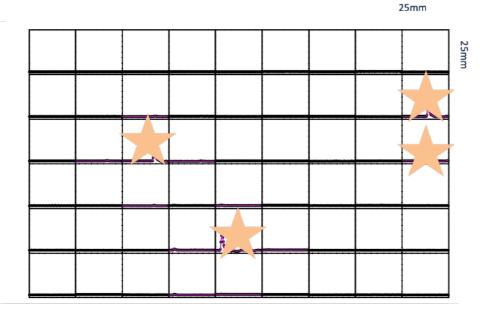
 $\omega_a$  blinded by both hardware and software

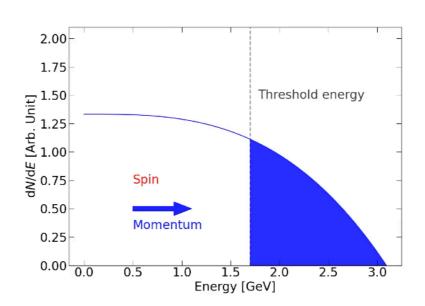


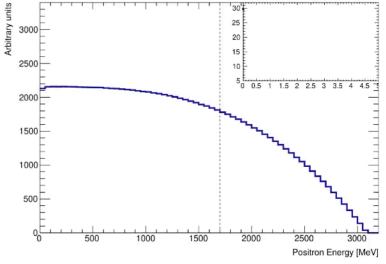




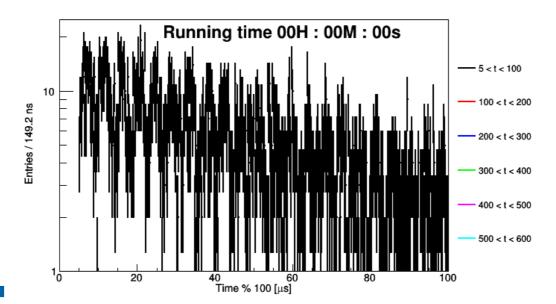
Energy distribution recorded in the calorimeter.

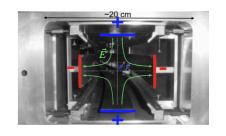




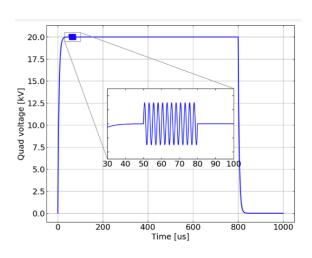


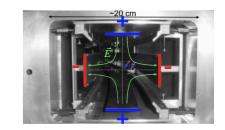
 $\omega_a$  blinded by both hardware and software

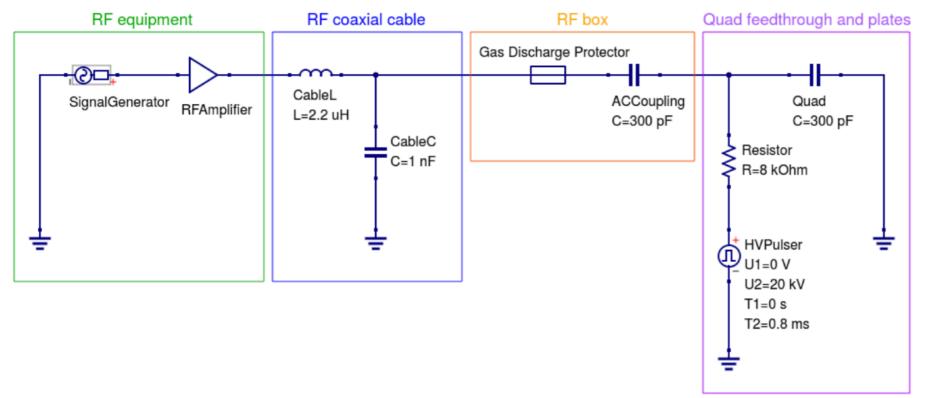




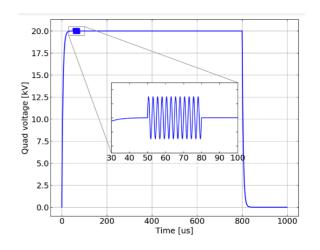
 RF can be applied through the quad plates

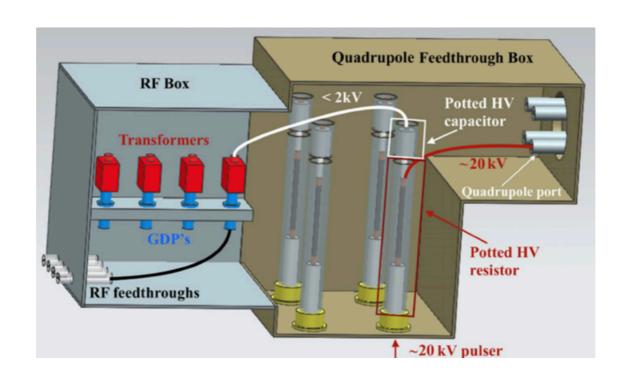


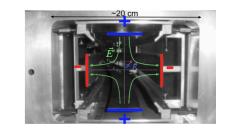


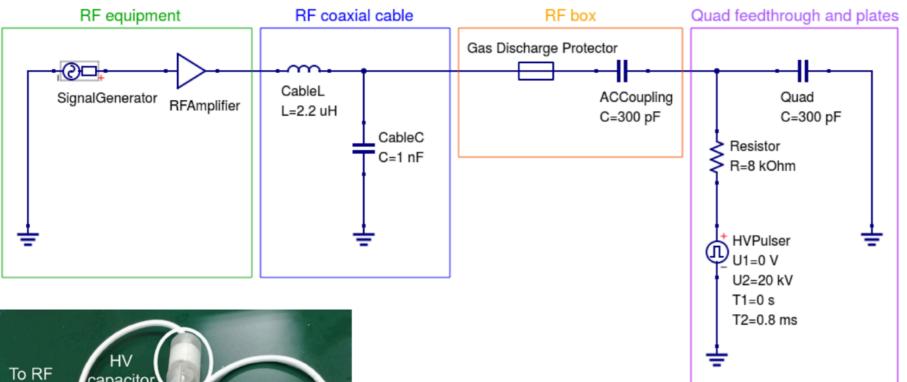


 RF can be applied through the quad plates

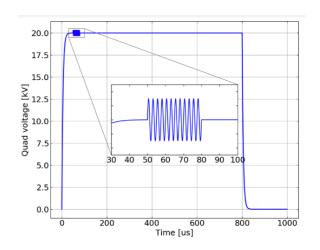


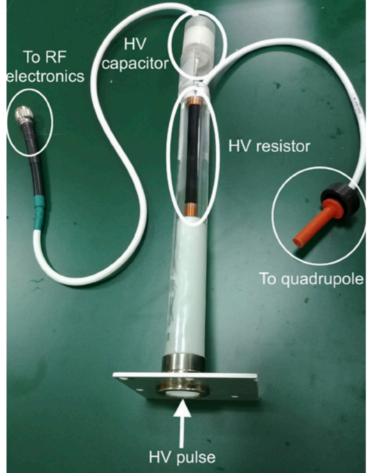


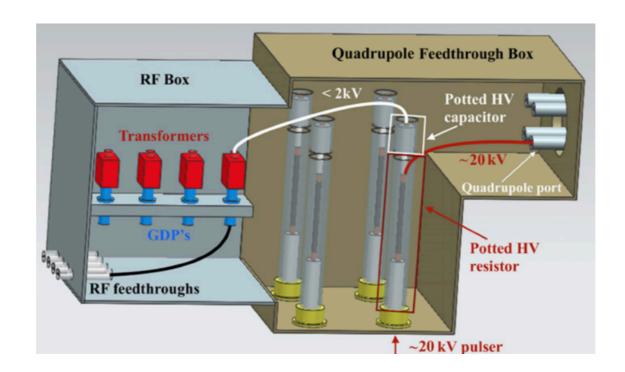




 RF can be applied through the quad plates

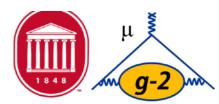






### What we gain from negative muon running

### CPTLV: $\mu^+/\mu^- \omega_a$ Difference



- Also note that Muon g-2 with  $\mu^-$  essentially gives you <u>3</u> new experimental results here, not just 1!
- What we have:
  - BNL( $\chi = 49.2$ )  $\mu^+/\mu^-$  (700 ppb) & CERN ( $\chi = 43.8$ )  $\mu^+/\mu^-$  (7300 ppb)  $\Delta cos \chi = 0.07$
- What we would get:
  - FNAL ( $\chi = 48.2$ )  $\mu^+$  (140 ppb) & BNL  $\mu^-$  (700 ppb)
  - FNAL  $\mu^{-}$  (280 ppb) & BNL  $\mu^{+}$  (700 ppb)
  - FNAL  $\mu$  (280 ppb) & J-PARC ( $\chi = 53.5$ )  $\mu$  (450 ppb) Dominates
- Potentially about 15x improvement

$$\Delta cos \chi = 0.01$$

$$\Delta cos \chi = 0.07$$



### **J-PARC Comparison**

# Comparison of Experiment Parameters

Table 1. Comparison of BNL-E821, FNAL-E989, and our experiment.

	BNL-E821	Fermilab-E989	Our experiment J-PARC E34	
Muon momentum	$3.09~{ m GeV}/c$		$300~{ m MeV}/c$	
Lorentz γ Radiu	Radius of cyclotron motion: 7.1 m		Radius of cyclotron	
Dalamination			50% motion: 333 mm	
Storage field			B = 3.0  T	
Focusing field	Electric quadrupole		Very weak magnetic	
Cyclotron period	149 1	ıs	7.4 ns	
Spin precession period	4.37	us	$2.11 \mu \mathrm{s}$	
Number of detected $e^+$	$5.0 \times 10^9$	$1.6 \times 10^{11}$	$5.7 \times 10^{11}$	
Number of detected e-	$3.6 \times 10^9$	_	_	
$a_{\mu}$ precision (stat.)	460 ppb	100 ppb	450 ppb	
(syst.)	280 ppb	100 ppb	<70 ppb	
EDM precision (stat.)	$0.2 \times 10^{-19} e \cdot cm$	_	$1.5 \times 10^{-21} e \cdot cm$	
(syst.)	$0.9 \times 10^{-19} e \cdot \text{cm}$	-	$0.36 \times 10^{-21} e \cdot \text{cm}$	

PTEP 2019 (2019), 053C02